Unresolved Anomalies and Tensions in the Standard Cosmological Model

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Eleonora Di Valentino Royal Society Dorothy Hodgkin Research Fellow School of Mathematics and Statistics University of Sheffield (UK)



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The **ACDM** model

Out of various cosmological models proposed in literature, the Lambda cold dark matter (ACDM) scenario has been chosen as the standard model for its simplicity and ability to accurately describe a wide range of astrophysical and cosmological observations.

However, ACDM still has many unknown areas and lacks the ability to explain fundamental concepts related to the structure and evolution of the universe. These concepts are based on three unknown ingredients that are not supported by theoretical first principles or laboratory experiments but are instead inferred from cosmological and astrophysical observations.

The three unknown ingredients are: inflation, dark matter (DM), and dark energy (DE). In ΛCDM, inflation is given by a single, slow-rolling scalar field; DM is assumed to interact only through gravity, be cold and pressureless, and lack direct evidence of its existence; DE is represented by the cosmological constant term Λ, without any strong physical explanation.

The **ACDM** model

Despite its theoretical shortcomings, ACDM remains the preferred model due to its ability to accurately describe observed phenomena.
However, the ACDM model with its six parameters is not based on deep-rooted physical principles and should be considered, at best,
an approximation of an underlying physical theory that remains undiscovered.

Hence, as observations become more numerous and accurate, deviations from the ΛCDM model are expected to be detected.
And in fact, discrepancies in important cosmological parameters, such as H0, have already arisen in various observations with different statistical significance.

While some of these tensions may have a systematic origin, their recurrence across multiple probes suggests that there may be flaws in the standard cosmological scenario, and that new physics may be necessary to explain these observational shortcomings.

Therefore, the persistence of these tensions could indicate the failure of the canonical ACDM model.

H0 tension

The H0 tension is the most statistically significant, long-lasting and widely persisting disagreement we have currently in cosmology.













Distance Ladder





Search...

Help | Advanced

Astrophysics > Cosmology and Nongalactic Astrophysics

[Submitted on 31 May 2023]

Leveraging SN Ia spectroscopic similarity to improve the measurement of ${\cal H}_0$

Yukei S. Murakami, Adam G. Riess, Benjamin E. Stahl, W. D'Arcy Kenworthy, Dahne-More A. Pluck, Antonella Macoretta, Dillon Brout, David O. Jones, Dan M. Scolnic, Alexei V. Filippenko

Recent studies suggest spectroscopic differences explain a fraction of the variation in Type Ia supernova (SN Ia) luminosities after light-curve/color standardization. In this work, (i) we empirically characterize the variations of standardized SN Ia luminosities, and (ii) we use a spectroscopically inferred parameter, SIP, to improve the precision of SNe Ia along the distance ladder and the determination of the Hubble constant (H_0). First, we show that the \texttt{Pantheon+} covariance model modestly overestimates the uncertainty of standardized magnitudes by ~ 7%, in the parameter space used by the SH0ES Team to measure H_0 ; accounting for this alone yields $H_0 = 73.01 \pm 0.92$ km s⁻¹ Mpc⁻¹. Furthermore, accounting for spectroscopic similarity between SNe~Ia on the distance ladder reduces their relative scatter to ~ 0.12 mag per object (compared to ~ 0.14 mag previously). Combining these two findings in the model of SN covariance, we find an overall 14% reduction (to ± 0.85 km s⁻¹ Mpc⁻¹) of the uncertainty in the Hubble constant and a modest increase in its value. Including a budget for systematic uncertainties itemized by Riess et al. (2022a), we report an updated local Hubble constant with ~ 1.2% uncertainty, $H_0 = 73.29 \pm 0.90$ km s⁻¹ Mpc⁻¹. We conclude that spectroscopic differences among photometrically standardized SNe Ia do not explain the ``Hubble tension." Rather, accounting for such differences its significance, as the discrepancy against Λ CDM calibrated by the *Planck* 2018 measurement rises to 5.7 σ .

The Planck estimate assuming a "vanilla" Λ CDM cosmological model: $H0 = 67.36 \pm 0.54 \text{ km/s/Mpc}$ Planck 2018, Astron.Astrophys. 641 (2020) A6 3σ -5σ -3 a -2 σ 2σ 5σ -4 σ -1 σ mean 1σ 4σ Derisity 50 Planck 2018 Baseline samples samples 0.0 73.04 67.472.03 10^{1} 74.06Density (log scale) 10^{-2} 10^{-2} 71.0475.10 70.07 76.1677.25 69.1068.1478.36 10^{-5} 0.99 1.04 1.06 68 70 76 66 72 74 78 H_0 (km/s/Mpc)

The latest local measurements obtained by the SH0ES collaboration

H0 = 73.04 ± 1.04 km/s/Mpc Riess et al. arXiv:2112.04510



The Universe originates from a hot Big Bang.

The primordial plasma in thermodynamic equilibrium cools with the expansion of the Universe. It goes through the phase of recombination, where electrons and protons combine into hydrogen atoms, and decoupling, where the Universe becomes transparent to the motion of photons.

The Cosmic Microwave Background (CMB) is the radiation from recombination, emitted about 13 billion years ago, just 400,000 years after the Big Bang.

The CMB provides an unparalleled probe of the early Universe and today it is a black body a temperature T=2.726K. T= $_{12}$



From the map of the CMB anisotropies we can extract the temperature angular power spectrum.





Cosmological parameters: ($\Omega_b h^2$, $\Omega_m h^2$, H0, n_s , τ , As)



We choose a set of cosmological parameters that describes our theoretical model and compute the angular power spectra. Because of the correlations present between the parameters, variation of different quantities can produce similar effects on the CMB.



Lemos & Shah, arXiv:2307.13083

Cosmological parameters: (Ω_bh^2 , Ω_mh^2 , H0, n_s, τ , As)



We compare the angular power spectra we computed with the data and, using a bayesian analysis, we get a combination of cosmological parameter values in agreement with these.



Planck 2018, Astron.Astrophys. 641 (2020) A6



Parameter constraints

We can extract 4 independent angular spectra from the CMB:

- Temperature
- Cross Temperature Polarization E
- Polarization type E (density fluctuations)
- Polarization type B (gravitational waves)



Borstnik et al., hep-ph/0401043

Planck satellite experiment



Planck 2018, Aghanim et al., arXiv:1807.06209 [astro-ph.CO]

Planck satellite experiment



The theoretical spectra in light blues are computed from the best-fit base-LCDM theoretical spectrum fit to the Planck TT,TE,EE+lowE+lensing likelihood.

Residuals with respect to this theoretical model are shown in the lower panel in each plot.



Polarization spectra

Planck 2018, Aghanim et al., arXiv:1807.06209 [astro-ph.CO]

Planck satellite experiment



Planck 2018, Astron.Astrophys. 641 (2020) A6

Constraints on parameters of the LCDM model from the separate Planck EE, TE, and TT high-I spectra combined with low-I polarization (lowE), and, in the case of EE also with BAO, compared to the joint result using Planck TT,TE,EE+lowE.

Parameter	TT+lowE 68% limits	TE+lowE 68% limits	EE+lowE 68% limits	TT,TE,EE+lowE 68% limits	TT,TE,EE+lowE+lensing 68% limits	TT,TE,EE+lowE+lensing+BAO 68% limits
$\overline{\Omega_{\rm b}h^2\ldots\ldots\ldots\ldots}$	0.02212 ± 0.00022	0.02249 ± 0.00025	0.0240 ± 0.0012	0.02236 ± 0.00015	0.02237 ± 0.00015	0.02242 ± 0.00014
$\Omega_{\rm c} h^2$	0.1206 ± 0.0021	0.1177 ± 0.0020	0.1158 ± 0.0046	0.1202 ± 0.0014	0.1200 ± 0.0012	0.11933 ± 0.00091
100 <i>θ</i> _{MC}	1.04077 ± 0.00047	1.04139 ± 0.00049	1.03999 ± 0.00089	1.04090 ± 0.00031	1.04092 ± 0.00031	1.04101 ± 0.00029
τ	0.0522 ± 0.0080	0.0496 ± 0.0085	0.0527 ± 0.0090	$0.0544^{+0.0070}_{-0.0081}$	0.0544 ± 0.0073	0.0561 ± 0.0071
$\ln(10^{10}A_s)$	3.040 ± 0.016	$3.018^{+0.020}_{-0.018}$	3.052 ± 0.022	3.045 ± 0.016	3.044 ± 0.014	3.047 ± 0.014
<i>n</i> _s	0.9626 ± 0.0057	0.967 ± 0.011	0.980 ± 0.015	0.9649 ± 0.0044	0.9649 ± 0.0042	0.9665 ± 0.0038
$H_0 [\mathrm{kms^{-1}Mpc^{-1}}]$	66.88 ± 0.92	68.44 ± 0.91	69.9 ± 2.7	67.27 ± 0.60	67.36 ± 0.54	67.66 ± 0.42
Ω_{Λ}	0.679 ± 0.013	0.699 ± 0.012	$0.711^{+0.033}_{-0.026}$	0.6834 ± 0.0084	0.6847 ± 0.0073	0.6889 ± 0.0056
$\Omega_m \ . \ . \ . \ . \ . \ . \ . \ . \ . \ $	0.321 ± 0.013	0.301 ± 0.012	$0.289^{+0.026}_{-0.033}$	0.3166 ± 0.0084	0.3153 ± 0.0073	0.3111 ± 0.0056
$\Omega_{\rm m}h^2$	0.1434 ± 0.0020	0.1408 ± 0.0019	$0.1404^{+0.0034}_{-0.0039}$	0.1432 ± 0.0013	0.1430 ± 0.0011	0.14240 ± 0.00087
$\Omega_{\rm m}h^3$	0.09589 ± 0.00046	0.09635 ± 0.00051	$0.0981^{+0.0016}_{-0.0018}$	0.09633 ± 0.00029	0.09633 ± 0.00030	0.09635 ± 0.00030
<i>σ</i> ₈	0.8118 ± 0.0089	0.793 ± 0.011	0.796 ± 0.018	0.8120 ± 0.0073	0.8111 ± 0.0060	0.8102 ± 0.0060
$S_8\equiv\sigma_8(\Omega_{\rm m}/0.3)^{0.5}$.	0.840 ± 0.024	0.794 ± 0.024	$0.781^{+0.052}_{-0.060}$	0.834 ± 0.016	0.832 ± 0.013	0.825 ± 0.011

Planck 2018, Astron.Astrophys. 641 (2020) A6

2018 Planck results are a wonderful confirmation of the flat standard ΛCDM cosmological model, but are **model dependent**!

- The cosmological constraints are obtained assuming a cosmological model.
- The results are affected by the degeneracy between the parameters that induce similar effects on the observables. 20



Are there other H0 estimates?



Hubble constant measurements made by different astronomical missions and groups over the years.

The red vertical band corresponds to the H0 value from SH0ES Team and the green vertical band corresponds to the H0 value as reported by Planck 2018 team within a ACDM scenario.

On the same side of Planck, i.e. preferring smaller values of H₀ we have:



Ground based CMB telescope



ACT-DR4 2020, JCAP 12 (2020) 047

ACT-DR4: H0 = 67.9 \pm 1.5 km/s/Mpc in Λ CDM

ACT-DR4 + WMAP: H0 = $67.6 \pm 1.1 \text{ km/s/Mpc}$ in ΛCDM

 ΛCDM - dependent

CMB Polarization Measurements with SPTpol

On the same side of Planck, i.e. preferring smaller values of H0 we have:

Ground based CMB telescope

Nicholas Harrington UC Berkeley



SPT-3G collaboration, arXiv:2212.05642

SPT-3G TT/TE/EE: H0 = 68.3 ± 1.5 km/s/Mpc in Λ CDM

 ΛCDM - dependent



On the same side of Planck, i.e. preferring smaller values of H₀ we have:

Spectroscopic Surveys BAO and Full Shape from BOSS and eBOSS

Results shown in blue include a BBN prior on ωb , in green use an ωb prior from *Planck*, in red are combined with the full *Planck* dataset.

 $_{\Lambda}CDM$ - dependent

Ivanov and Philcox, arXiv:2305.07977





Freedman, arXiv:2106.15656









H0 = 76.00 ± 2.55 km/s/Mpc Kourkchi et al. arXiv:2004.14499

H0 = 75.10 ± 2.75 km/s/Mpc Schombert et al. arXiv:2006.08615

Tully-Fisher Relation (based on the correlation between the rotation rate of spiral galaxies and their absolute luminosity, and using as calibrators Cepheids and TRGB)







Freedman, *Astrophys.J.* 919 (2021) 1, 16

In the past the tension was within the same types of measurements and at the same redshifts and thus pointing directly to systematics.


It is difficult to imagine a single systematic error that would consistently explain the discrepancies observed in the diverse range of phenomena that we have encountered earlier, thereby resolving the Hubble constant tension.

Since this tension persists in the 5 - 6.3σ range

(Riess, Nature Reviews Physics (2019); Di Valentino, MNRAS 502 (2021) 2, 2065-2073; Di Valentino, Universe 2022, 8(8), 399) even after eliminating the measurements of any individual type of object, team, or calibration, it is challenging to identify a single error that could account for it. While multiple independent systematic errors could offer more flexibility in resolving the tension, they are less likely to occur.

Given that the indirect constraints are model-dependent, we can explore the possibility of expanding the cosmological scenario and examining which extensions can resolve the discrepancies between the various cosmological probes.

Let's modify the ACDM model with a few example...

(Di Valentino et al. Class. Quant. Grav. 38 (2021) 15, 153001 and Abdalla et al., JHEAp 34 (2022) 49-211)

The Neutrino effective number

We can consider modifications in the dark matter sector.

A classical extension is the effective number of relativistic degrees of freedom, i.e. additional relativistic matter at recombination, corresponding to a modification of the expansion history of the universe at early times.

The Neutrino effective number

The expected value is Neff = 3.044, if we assume standard electroweak interactions and three active massless neutrinos. If we measure a Neff > 3.044, we are in presence of extra radiation.

If we vary Neff, at 68% cl H0 is equal to 66.4 ± 1.4 km/s/Mpc, and the tension with SH0ES is still 3.9σ .

 $N_{\rm eff} = 2.92^{+0.36}_{-0.37}$ (95%, *Planck* TT, TE, EE+lowE),



Planck 2018, Astron.Astrophys. 641 (2020) A6

The Dark energy equation of state

For example, we can consider modifications in the dark energy sector.

A classical extension is a varying dark energy equation of state, that is a modification of the expansion history of the universe at late times.

The Dark energy equation of state

If we change the cosmological constant with a Dark Energy with equation of state w, we are changing the expansion rate of the Universe:

$$H^{2} = \left(\frac{\dot{a}}{a}\right)^{2} = H_{0}^{2} \left(\frac{\Omega_{r}}{a^{4}} + \frac{\Omega_{m}}{a^{3}} + \frac{\Omega_{k}}{a^{2}} + \Omega_{\Lambda}\right)$$

$$H^{2} = H_{0}^{2} \left[\Omega_{m} (1+z)^{3} + \Omega_{r} (1+z)^{4} + \Omega_{de} (1+z)^{3(1+w)} + \Omega_{k} (1+z)^{2} \right]$$

w introduces a geometrical degeneracy with the Hubble constant that is almost unconstrained using the CMB data only, resulting in agreement with SH0ES.

We have in 2018 w = $-1.58^{+0.52}_{-0.41}$ with H0 > 69.9 km/s/Mpc at 95% c.l.

Planck data prefer a phantom dark energy, with an energy component with w < -1, for which the density increases with time in an expanding universe that will end in a Big Rip. A phantom dark energy violates the energy condition $\rho \ge |p|$, that means that the matter could move faster than light and a comoving observer measure a negative energy density, and the Hamiltonian could have vacuum instabilities due to a negative kinetic energy.

Formally successful models in solving H0

tension $\leq 1\sigma$ "Excellent models"	tension $\leq 2\sigma$ "Good models"	tension $\leq 3\sigma$ "Promising models"
Dark energy in extended parameter spaces [289]	Early Dark Energy [235]	Early Dark Energy [229]
Dynamical Dark Energy [309]	Phantom Dark Energy [11]	Decaying Warm DM [474]
Metastable Dark Energy [314]	Dynamical Dark Energy [11, 281, 309]	Neutrino-DM Interaction [506]
PEDE [392, 394]	GEDE [397]	Interacting dark radiation [517]
Elaborated Vacuum Metamorphosis [400–402]	Vacuum Metamorphosis [402]	Self-Interacting Neutrinos [700, 701]
IDE [314, 636, 637, 639, 652, 657, 661–663]	IDE [314,653,656,661,663,670]	IDE [656]
Self-interacting sterile neutrinos [711]	Critically Emergent Dark Energy [997]	Unified Cosmologies [747]
Generalized Chaplygin gas model [744]	$f(\mathcal{T})$ gravity [814]	Scalar-tensor gravity [856]
Galileon gravity [876, 882]	Über-gravity [59]	Modified recombination [986]
Power Law Inflation [966]	Reconstructed PPS [978]	Super ΛCDM [1007]
$f(\mathcal{T})$ [818]		Coupled Dark Energy [650]

Table B1. Models solving the H_0 tension with R20 within the 1σ , 2σ and 3σ Planck only confidence levels considering the *Planck* dataset only.

Di Valentino et al., Class.Quant.Grav. (2021), arXiv:2103.01183 [astro-ph.CO]

The state of the Dark energy equation of state

Dataset combination	$oldsymbol{w}$	$H_0[{ m km/s/Mpc}]$	
CMB	$-1.57^{+0.16}_{-0.36} \ (-1.57^{+0.53}_{-0.42})$	> 82.4 (> 69.3)	
CMB+BAO	$-1.039 \pm 0.059 \ (-1.04^{+0.11}_{-0.12})$	$68.6 \pm 1.5 (68.6^{+3.1}_{-2.8})$	
CMB+SN	$-0.976 \pm 0.029\;(-0.976 \substack{+0.055 \\ -0.056})$	$66.54 \pm 0.81 (66.5^{+1.6}_{-1.6})$	

Escamilla, Giarè, Di Valentino et al., arXiv: 2307.14802



FIG. 5. Best-fit predictions for (rescaled) distance-redshift relations from a wCDM fit to *Planck* CMB data alone (dashed curves) and the CMB+BAO dataset (solid curves). These predictions are presented for the three different types of distances probed by BAO measurements (rescaled as per the y label), each indicated by the colors reported in the legend. The error bars represent $\pm 1\sigma$ uncertainties. However, if BAO data are included, the wCDM model with w<-1 worsens considerably the fit of the BAO data because the best fit from Planck alone fails in recover the shape of H(z) at low redshifts. Therefore, when the CMB is combined with BAO data, the favoured model is again the LCDM one and the H0 tension is restored.

Complication: the sound horizon problem

What about BAO+Pantheon?

BAO+Pantheon measurements constrain the product of H0 and the sound horizon r_s.

In order to have a higher H0 value in agreement with SH0ES, we need r_s near 137 Mpc. However, Planck by assuming Λ CDM, prefers r_s near 147 Mpc. Therefore, a cosmological solution that can increase H0 and at the same time can lower the sound horizon inferred from CMB data is the most promising way to put in agreement all the measurements.



Knox and Millea, *Phys.Rev.D* 101 (2020) 4, 043533

Early vs late time solutions

Here we can see the comparison of the 2 σ credibility regions of the CMB constraints and the measurements from late-time observations (SN + BAO + H0LiCOW + SH0ES).

We see that the late time solutions, as wCDM, increase H0 because they decrease the expansion history at intermediate redshift, but leave r_s unaltered.



Arendse et al., Astron.Astrophys. 639 (2020) A57

Early vs late time solutions

Here we can see the comparison of the 2 σ credibility regions of the CMB constraints and the measurements from late-time observations (SN + BAO + H0LiCOW + SH0ES).

However, the early time solutions, as Neff or Early Dark Energy, move in the right direction both the parameters, but can't solve completely the H0 tension between Planck and SH0ES.



Arendse et al., Astron.Astrophys. 639 (2020) A57

Early Dark Energy

Early dark energy (EDE) scenario assumes that there is a new fundamental field that accelerates the cosmic expansion rate before recombination. This field contributes roughly 10-12% of the total energy density near the matter-radiation equality, but eventually dissipates like radiation or at a faster rate (depending on the shape of the potential). In order to have an effect on the sound horizon we should have H ~ T²/M_{pl}≈ m just before the recombination, so the mass of the scalar field should be m ≈ 10⁻²⁷ eV, similar to an axion particle:

$$V(\phi) = m^2 f^2 (1 - \cos(\phi/f))^n$$

At the minimum of the potential the field oscillates yielding to an effective equation of state

$$w_{\phi} = (n-1)/(n+1)$$

If we take n = 1 (the standard axion potential) then w_o = 0 near the potential minimum, and the EDE energy density redshifts as matter creating problems in the late-time cosmology, therefore it does not work phenomenologically.
For n = 2 instead it decays away like radiation (∝ a⁻⁴), and for n → ∞ like kinetic energy (∝ a⁻⁶). However, values n > 5 are disfavored.

Early Dark Energy

Constraints at 68% cl.

Constraints from <i>Planck</i> 2018 data only: TT+TE+EE				
Parameter	ΛCDM	EDE $(n = 3)$		
$\ln(10^{10}A_{\rm s})$	$3.044(3.055)\pm 0.016$	$3.051(3.056)\pm 0.017$		
$n_{ m s}$	$0.9645(0.9659)\pm 0.0043$	$0.9702(0.9769)^{+0.0071}_{-0.0069}$		
$100 heta_{ m s}$	$1.04185(1.04200) \pm 0.00029$	$ig 1.04164 (1.04168) \pm 0.00034$		
$ \Omega_{ m b}h^2$	$0.02235(0.02244) \pm 0.00015$	$ig 0.02250(0.02250)\pm 0.00020$		
$ \Omega_{ m c}h^2$	$0.1202(0.1201)\pm 0.0013$	$0.1234(0.1268)^{+0.0031}_{-0.0030}$		
$ au_{ m reio}$	$0.0541(0.0587)\pm 0.0076$	$0.0549(0.0539)\pm 0.0078$		
$\log_{10}(z_c)$	—	$3.66(3.75)^{+0.28}$		
$f_{ m EDE}$	_	< 0.087(0.068)		
$oldsymbol{ heta}_{i}$	_	> 0.30 (2.96)		
$H_0[{ m km/s/Mpc}]$	$67.29(67.44)\pm 0.59$	$68.29(69.13)^{+1.02}_{-1.00}$		
$\Omega_{ m m}$	$0.3162(0.3147)\pm 0.0083$	$0.3145(0.3138)\pm 0.0086$		
σ_8	$0.8114(0.8156)\pm 0.0073$	$0.8198(0.8280)^{+0.0109}_{-0.0107}$		
S_8	$0.8331(0.8355)\pm 0.0159$	$0.8393(0.8468)\pm 0.0173$		
$\log_{10}(f/{ m eV})$	_	$26.57(26.36)^{+0.39}_{-0.36}$		
$\log_{10}(m/{ m eV})$	_	$-26.94 (-26.90)^{+0.58}_{-0.53}$		

Hill et al. Phys.Rev.D 102 (2020) 4, 043507

Planck 2018 results shows no evidence for EDE and H0 is in agreement with the value obtained assuming ACDM.

Formally successful models in solving H0

tension $\leq 1\sigma$ "Excellent models"	tension $\leq 2\sigma$ "Good models"	tension $\leq 3\sigma$ "Promising models"
Early Dark Energy [228, 235, 240, 250]	Early Dark Energy [212, 229, 236, 263]	DE in extended parameter spaces [289]
Exponential Acoustic Dark Energy [259]	Rock 'n' Roll [242]	Dynamical Dark Energy [281, 309]
Phantom Crossing [315]	New Early Dark Energy [247]	Holographic Dark Energy [350]
Late Dark Energy Transition [317]	Acoustic Dark Energy [257]	Swampland Conjectures [370]
Metastable Dark Energy [314]	Dynamical Dark Energy [309]	MEDE [399]
PEDE [394]	Running vacuum model [332]	Coupled DM - Dark radiation [534]
Vacuum Metamorphosis [402]	Bulk viscous models [340, 341]	Decaying Ultralight Scalar [538]
Elaborated Vacuum Metamorphosis [401, 402]	Holographic Dark Energy [350]	$BD-\Lambda CDM$ [852]
Sterile Neutrinos [433]	Phantom Braneworld DE [378]	Metastable Dark Energy [314]
Decaying Dark Matter [481]	PEDE [391, 392]	Self-Interacting Neutrinos [700]
Neutrino-Majoron Interactions [509]	Elaborated Vacuum Metamorphosis [401]	Dark Neutrino Interactions [716]
IDE [637, 639, 657, 661]	IDE [659,670]	$IDE \ [634-636, 653, 656, 663, 669]$
DM - Photon Coupling [685]	Interacting Dark Radiation [517]	Scalar-tensor gravity [855,856]
$f(\mathcal{T})$ gravity theory [812]	Decaying Dark Matter [471, 474]	Galileon gravity [877,881]
BD- Λ CDM [851]	DM - Photon Coupling [686]	Nonlocal gravity [886]
Über-Gravity [59]	Self-interacting sterile neutrinos [711]	Modified recombination [986]
Galileon Gravity [875]	$f(\mathcal{T})$ gravity theory [817]	Effective Electron Rest Mass [989]
Unimodular Gravity [890]	Über-Gravity [871]	Super ACDM [1007]
Time Varying Electron Mass [990]	VCDM [893]	Axi-Higgs [991]
M CDM [995]	Primordial magnetic fields [992]	Self-Interacting Dark Matter [479]
Ginzburg-Landau theory [996]	Early modified gravity [859]	Primordial Black Holes [545]
Lorentzian Quintessential Inflation [979]	Bianchi type I spacetime [999]	
Holographic Dark Energy [351]	f(T) [818]	

Combination of datasets **Table B2.** Models solving the H_0 tension with R20 within 1σ , 2σ and 2σ *Planck* in combination with additional cosmological probes. datasets are discussed in the main text.

Di Valentino et al., Class.Quant.Grav. (2021), arXiv:2103.01183 [astro-ph.CO]

Additional complication: the early solutions proposed to alleviate the H0 tension increase the S8 tension!



$$S_8 \equiv \sigma_8 \sqrt{\Omega_m/0.3}$$

A tension on S8 is present between the Planck data in the ACDM scenario and the cosmic shear data.



KiDS-1000, Heymans et al., arXiv:2007.15632 [astro-ph.CO]

The S8 tension is present at 3.4σ between Planck assuming ΛCDM and KiDS+VIKING-450 and BOSS combined together, or 3.1σ with KiDS-1000.

 $S_8 = 0.834 \pm 0.016$ Planck 2018, Aghanim et al., arXiv:1807.06209 [astro-ph.CO]

> $S_8 = 0.728 \pm 0.045$ Troster et al., arXiv:1909.11006 [astro-ph.CO]

> > $S_8 = 0.766^{+0.020}_{-0.014}$

KiDS-1000, Heymans et al., arXiv:2007.15632 [astro-ph.CO]



DES-Y3, Amon et al., arXiv:2105.13543 [astro-ph.CO]

The S8 tension is present at 2.5σ between Planck assuming Λ CDM and DES-Y3.

 $S_8 = 0.834 \pm 0.016$

Planck 2018, Aghanim et al., arXiv:1807.06209 [astro-ph.CO]

 $S_8 = 0.776^{+0.017}_{-0.017}$

DES-Y3, Abbott et al., arXiv:2105.13549 [astro-ph.CO]

 $S_8 = 0.759^{+0.025}_{-0.025}$

DES-Y3 fiducial, Amon et al., arXiv:2105.13543 [astro-ph.CO]



HSC-Y3, Dalal et al., arXiv:2304.00701 [astro-ph.CO]

The S8 tension is present at about 2σ between Planck assuming ΛCDM and HSC-Y3.

 $S_8 = 0.834 \pm 0.016$

Planck 2018, Aghanim et al., arXiv:1807.06209 [astro-ph.CO]

 $S_8 = 0.776^{+0.032}_{-0.033}$

HSC-Y3, Dalal et al., arXiv:2304.00701 [astro-ph.CO]

Aghanim et al. (2020d)

Aghanim et al. (2020d)

Aiola et al. (2020)

• CMB Planck TT,TE,EE+lowE • CMB Planck TT,TE,EE+lowE+lensing • CMB ACT+WMAP



See Di Valentino et al. *Astropart.Phys.* 131 (2021) 102604 and Abdalla et al., arXiv:2203.06142 [astro-ph.CO] for a summary of the possible candidates proposed to solve the S8 tension.

Early solutions to the H0 tension

Actually, a dark energy model that merely changes the value of rd would not completely resolve the tension, since it will affect the inferred value of Ω m and transfer the tension to it.

This is a plot illustrating that achieving a full agreement between CMB, BAO and SH0ES through a reduction of rd requires a higher value of $\Omega_m h^2$.



Jedamzik et al., Commun.in Phys. 4 (2021) 123

Early solutions to the H0 tension

Model 2 is defined by the simultaneous fit to BAO and CMB acoustic peaks at $\Omega_m h^2 = 0.155$, while model 3 has $\Omega_m h^2 = 0.167$

The sound horizon problem should be considered not only in the plane H0–rd, but it should be extended to the parameters triplet H0–rd– Ω m.

The figure shows that when attempting to find a full resolution of the Hubble tension, with CMB, BAO and SH0ES in agreement with each other, one exacerbates the tension with DES, KiDS and HSC.



Jedamzik et al., Commun.in Phys. 4 (2021) 123

Successful models?

This is the density of the proposed cosmological models:

At the moment no specific proposal makes a strong case for being highly likely or far better than all others !!!



Di Valentino et al., Class.Quant.Grav. (2021), arXiv:2103.01183 [astro-ph.CO]

What about the interacting DM-DE models?

In the standard cosmological framework, DM and DE are described as separate fluids not sharing interactions beyond gravitational ones.

At the background level, the conservation equations for the pressureless DM and DE components can be decoupled into two separate equations with an inclusion of an arbitrary function, *Q*, known as the coupling or interacting function:

$$\dot{\rho}_c + 3\mathcal{H}\rho_c = Q,$$

$$\dot{\rho}_x + 3\mathcal{H}(1+w)\rho_x = -Q,$$

and we assume the phenomenological form for the interaction rate:

$$Q = \xi \mathcal{H} \rho_x$$

proportional to the dark energy density ρ_x and the conformal Hubble rate \mathcal{H} , via a negative dimensionless parameter ξ quantifying the strength of the coupling, to avoid early-time instabilities.

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In this scenario of IDE the tension on H0 between the Planck satellite and SH0ES is completely solved. The coupling could affect the value of the present matter energy density Ω_m . Therefore, if within an interacting model Ω_m is smaller (because for negative ξ the dark matter density will decay into the dark energy one), a larger value of H0 would be required in order to satisfy the peaks structure of CMB observations, which accurately determine the value of $\Omega_m h^2$.

	Parameter	Planck	Planck+R19
	$\Omega_{ m b}h^2$	0.02239 ± 0.00015	0.02239 ± 0.00015
	$\Omega_{ m c} h^2$	< 0.105	< 0.0615
	n_s	0.9655 ± 0.0043	0.9656 ± 0.0044
	$100\theta_{s}$	$1.0458\substack{+0.0033\\-0.0021}$	1.0470 ± 0.0015
	au	0.0541 ± 0.0076	0.0534 ± 0.0080
	ξ	$-0.54^{+0.12}_{-0.28}$	$-0.66\substack{+0.09\\-0.13}$
H_0 [$[{\rm kms^{-1}Mpc^{-1}}]$	$72.8^{+3.0}_{-1.5}$	$74.0^{+1.2}_{-1.0}$

TABLE I. Mean values with their 68% C.L. errors on selected cosmological parameters within the $\xi \Lambda \text{CDM}$ model, considering either the *Planck* 2018 legacy dataset alone, or the same dataset in combination with the *R19* Gaussian prior on H_0 based on the latest local distance measurement from *HST*. The quantity quoted in the case of $\Omega_c h^2$ is the 95% C.L. upper limit.

Therefore we can safely combine the two datasets together, and we obtain a nonzero dark matter-dark energy coupling ξ at more than FIVE standard deviations.



Di Valentino et al., Phys.Dark Univ. 30 (2020) 100666 64

Moreover, we find a shift of the clustering parameter σ_8 towards a higher value, compensated by a lowering of the matter density Ω_m , both with relaxed error bars. The reason is that once a coupling is switched on and Ω_m becomes smaller, the clustering parameter σ_8 must be larger to have a proper normalization of the (lensing and clustering) power spectra.

This model can therefore significantly reduce the significance of the S8 tension (See also Lucca, *Phys.Dark Univ.* 34 (2021) 100899)



Di Valentino et al., Phys.Dark Univ. 30 (2020) 100666 65

Bayes factor

Anyway it is clearly interesting to quantify the better accordance of a model with the data respect to another by using the marginal likelihood also known as the Bayesian evidence.

The Bayesian evidence weights the simplicity of the model with the improvement of the fit of the data. In other words, because of the Occam's razor principle, models with additional parameters are penalised, if don't improve significantly the fit.

Given two competing models M₀ and M₁ it is useful to consider the ratio of the likelihood probability (the Bayes factor):

 $ln\mathcal{B} = p(\boldsymbol{x}|M_0)/p(\boldsymbol{x}|M_1)$

According to the revised Jeffrey's scale by Kass and Raftery 1995, the evidence for M_0 (against M_1) is considered as "weak" if | InB | > 1.0, "moderate" if | InB | > 2.5, and "strong" if | InB | > 5.0.

Parameter	Planck	Planck+R19
$\Omega_{ m b}h^2$	0.02239 ± 0.00015	0.02239 ± 0.00015
$\Omega_{ m c}h^2$	< 0.105	< 0.0615
n_s	0.9655 ± 0.0043	0.9656 ± 0.0044
$100\theta_{s}$	$1.0458\substack{+0.0033\\-0.0021}$	1.0470 ± 0.0015
au	0.0541 ± 0.0076	0.0534 ± 0.0080
ξ	$-0.54^{+0.12}_{-0.28}$	$-0.66^{+0.09}_{-0.13}$
$H_0 [{\rm km s^{-1} Mpc^{-1}}]$	$72.8^{+3.0}_{-1.5}$	$74.0^{+1.2}_{-1.0}$

TABLE I. Mean values with their 68% C.L. errors on selected cosmological parameters within the $\xi \Lambda \text{CDM}$ model, considering either the *Planck* 2018 legacy dataset alone, or the same dataset in combination with the *R19* Gaussian prior on H_0 based on the latest local distance measurement from *HST*. The quantity quoted in the case of $\Omega_c h^2$ is the 95% C.L. upper limit.

Di Valentino et al., Phys.Dark Univ. 30 (2020) 100666

Computing the Bayes factor for the IDE model with respect to ACDM for the Planck dataset we find InB = 1.2, i.e. a weak evidence for the IDE model. If we consider Planck + SH0ES we find the extremely high value InB=10.0, indicating a strong evidence for the IDE model.

fake IDE detection

Parameters	Fiducial model	Planck	Planck+BAO	PICO	PRISM
$\Omega_b h^2 \ \Omega_c h^2$	0.02236 0.1202	$\begin{array}{c} 0.02238 \pm 0.00015 \\ 0.056 \substack{+0.025 \\ -0.047 \\ \end{array}$	$\begin{array}{c} 0.02230 \pm 0.00014 \\ 0.101 \substack{+0.019 \\ -0.006} \end{array}$	$\begin{array}{c} 0.022364 \pm 0.000029 \\ 0.100^{+0.019}_{-0.008} \end{array}$	$\begin{array}{c} 0.022361 \pm 0.000019 \\ 0.103 +0.016 \\ -0.007 \\ -0.007 \\ -0.007 \\ +0.016 \\ -0.007 \\ +0.016 \\ -0.007 \\ +0.000 \\ -0.007 \\ +0.000 \\ -0.007 \\ +0.000 \\ -0.000 \\ +0.000 \\ -0.000 \\ +0.000 \\ -0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0.000 \\ +0$
$100 \theta_{MC} \ au$	1.04090 0.0544	$\frac{1.0451^{+0.0021}_{-0.0032}}{0.0528^{+0.010}_{-0.009}}$	$\begin{array}{c} 1.0419\substack{+0.0005\\-0.0011}\\ 0.0517\pm0.0098\end{array}$	$\frac{1.04206^{+0.0005}_{-0.0011}}{0.0543^{+0.0016}_{-0.0019}}$	$\frac{1.04191^{+0.00042}_{-0.00094}}{0.0542^{+0.0017}_{-0.0019}}$
n_s $\ln(10^{10} \Lambda)$	0.9649	0.9652 ± 0.0041	0.9624 ± 0.0036	0.9571 ± 0.0014	0.9657 ± 0.0012
$f(10^{-s}A_s)$	0	$-0.48^{+0.16}_{-0.30}$	> -0.223	> -0.220	> −0.195

Di Valentino & Mena, Mon.Not.Roy.Astron.Soc. 500 (2020) 1, L22-L26, arXiv:2009.12620

For a mock Planck-like experiment,

due to the strong correlation present between the standard and the exotic physics parameters, there is a dangerous detection at more than 3σ for a coupling between dark matter and dark energy different from zero, even if the fiducial model has $\xi = 0$:

 $-0.85 < \xi < -0.02$ at 99% CL



Mock experiments

fake IDE detection

Parameters	Fiducial model	Planck	Planck+BAO	PICO	PRISM
$\Omega_b h^2 \ \Omega_c h^2$	0.02236 0.1202	$\begin{array}{c} 0.02238 \pm 0.00015 \\ 0.056^{+0.025}_{-0.047} \end{array}$	$\begin{array}{c} 0.02230 \pm 0.00014 \\ 0.101 \substack{+0.019 \\ -0.006} \end{array}$	$\begin{array}{r} 0.022364 \pm 0.000029 \\ 0.100^{+0.019}_{-0.008} \end{array}$	$\begin{array}{r} 0.022361 \pm 0.000019 \\ 0.103 \substack{+0.016 \\ -0.007} \end{array}$
$100 \theta_{MC}$	1.04090	$1.0451_{-0.0032}^{+0.0021}$	$1.0419_{-0.0011}^{+0.0005}$	$1.04206_{-0.0011}^{+0.0005}$	$1.04191^{+0.00042}_{-0.00094}$
au	0.0544	$0.0528_{-0.009}^{+0.010}$	0.0517 ± 0.0098	$0.0543^{+0.0016}_{-0.0019}$	$0.0542^{+0.0017}_{-0.0019}$
n_s	0.9649	0.9652 ± 0.0041	0.9624 ± 0.0036	0.9571 ± 0.0014	0.9657 ± 0.0012
$\ln(10^{10}A_s)$	3.045	$3.041^{+0.020}_{-0.018}$	3.042 ± 0.019	$3.0436^{+0.0030}_{-0.0034}$	3.0435 ± 0.0032
ξ	0	$-0.43^{+0.16}_{-0.30}$	> -0.223	> -0.220	> -0.195

Di Valentino & Mena, Mon.Not.Roy.Astron.Soc. 500 (2020) 1, L22-L26, arXiv:2009.12620

The inclusion of mock BAO data, a mock dataset built using the same fiducial cosmological model than that of the CMB, helps in breaking the degeneracy, providing a lower limit for the coupling ξ in perfect agreement with zero. $1.00 - \frac{1.00}{0.75} - \frac{0.75}{0.50} - \frac{0.25}{0.25} - \frac{0.25}{0.00} - \frac{0.00}{-1.0 - 0.8 - 0.6 - 0.4 - 0.2 0.0} - \frac{\xi}{\xi}$

Mock experiments

Constraints at 68% cl.

Parameter	CMB+BAO	CMB+FS	CMB+BAO+FS
ω_c	$0.094\substack{+0.022\\-0.010}$	$0.101\substack{+0.015\\-0.009}$	$0.115\substack{+0.005\\-0.001}$
ξ	$-0.22^{+0.18}_{-0.09}$ [> -0.4	[48] > -0.35	> -0.12
$\left H_0 \mathrm{[km/s/Mpc]} ight $	$69.55\substack{+0.98\\-1.60}$	$69.04\substack{+0.84 \\ -1.10}$	$68.02\substack{+0.49\\-0.60}$
Ω_m	$0.243\substack{+0.054\\-0.030}$	$0.261\substack{+0.038\\-0.025}$	$0.299\substack{+0.015\\-0.007}$

Nunes, Vagnozzi, Kumar, Di Valentino, and Mena, Phys. Rev. D 105 (2022) 12, 123506

The addition of low-redshift measurements, as BAO data, still hints to the presence of a coupling, albeit at a lower statistical significance. Also for this data sets the Hubble constant value is larger than that obtained in the case of a pure Λ CDM scenario, enough to bring the H0 tension at 2.1 σ with SH0ES.

Baryon Acoustic Oscillations

BAO is formed in the early universe, when baryons are strongly coupled to photons, and the gravitational collapse due to the CDM is counterbalanced by the radiation pressure. Sound waves that propagate in the early universe imprint a characteristic scale on the CMB. Since the scale of these oscillations can be measured at recombination, BAO is considered a "standard ruler". These fluctuations have evolved and we can observe BAO at low redshifts in the distribution of galaxies.

Since the data reduction process leading to these measurements involves making certain assumptions about the fiducial cosmology, this makes BAO measurements dependent on the cosmological model being used.



Baryon Acoustic Oscillations

In other words, the tension between Planck+BAO and SH0ES could be due to a statistical fluctuation in this case.

Actually, BAO data are extracted under the assumption of ACDM, and the modified scenario of interacting dark energy could affect the result.

In fact, the full procedure which leads to the BAO datasets carried out by the different collaborations might be not necessarily valid in extended DE models with important perturbations in the non-linear scales.

BAO datasets (both the pre- and post- reconstruction measurements) might need to be revised in a non-trivial manner when applied to constrain more exotic dark energy cosmologies.
Baryon Acoustic Oscillations

The problem is that for 3D BAO data one needs to reconstruct the comoving distance and this is done assuming a fiducial model.

We can try to see what happens using 2D BAO measurements, that are less model dependent because they are obtained working on spherical shells with redshift thickness Δz and only considering their angular distribution.

Parameter	Planck	+ lensing	Planck + BAO	+ lensing	Pla	anck + BAOtr	+ lensing	$\mathbf{Planck} + \mathbf{BAOtr} + H_0$	+ lensing
H ₀ [Km/s/Mpc]	67.32 ± 0.62	67.32 ± 0.53	67.65 ± 0.44	67.60 ± 0.43		69.01 ± 0.51	68.85 ± 0.55	69.88 ± 0.48	69.65 ± 0.44
S_8	0.832 ± 0.016	0.834 ± 0.013	0.825 ± 0.012	0.827 ± 0.011		0.794 ± 0.013	0.802 ± 0.012	0.774 ± 0.013	$0.7871\substack{+0.0095\\-0.011}$
$r_s \; [{ m Mpc}]$	$\boxed{147.06\pm0.30}$	147.04 ± 0.27	$147.21\substack{+0.23 \\ -0.26}$	147.13 ± 0.23		147.75 ± 0.26	147.64 ± 0.26	148.06 ± 0.25	147.91 ± 0.24

A comparison between the **3D BAO data**, model dependent and obtained assuming ΛCDM, and the **2D BAO measurements**, less model dependent, shows almost the same results for the ΛCDM scenario.

Bernui, Di Valentino, Giarè, Kumar, and Nunes, Phys. Rev. D 107 (2023) 10, 103531

Parame	eter	Planck		Planck + BAO		Planck + BAOtr		Planck + BAOtr + H_0	
			+ lensing		+ lensing		+ lensing		+ lensing
\$ 68%	% CL	$-0.43\substack{+0.28\\-0.21}$	$-0.40\substack{+0.23\\-0.20}$	> -0.207	> -0.210	$-0.683\substack{+0.088\\-0.11}$	$-0.683\substack{+0.087\\-0.12}$	-0.58 ± 0.11	-0.53 ± 0.11
95%	% CL	(> -0.775)	$(-0.40\substack{+0.40\\-0.32})$	(> -0.389)	(> -0.411)	$(-0.68\substack{+0.21\\-0.19})$	$(-0.68\substack{+0.23\\-0.20})$	$(-0.58\substack{+0.22\-0.21})$	$(-0.53\substack{+0.19\\-0.20})$
99%	% CL	[> -0.819]	[> -0.743]	[> -0.486]	[> -0.527]	$[-0.68\substack{+0.29\\-0.23}]$	$[-0.68\substack{+0.37\\-0.27}]$	$[-0.58\substack{+0.31\\-0.29}]$	$[-0.53\substack{+0.39\\-0.25}]$
H ₀ [Kn	n/s/Mpc]	$71.7^{+2.3}_{-2.7}$	71.6 ± 2.1	$68.93\substack{+0.79\\-1.2}$	$69.08\substack{+0.74 \\ -1.3}$	$75.2^{+1.2}_{-0.75}$	$75.3^{+1.3}_{-0.75}$	73.99 ± 0.88	$73.45\substack{+0.71 \\ -0.59}$
S_8		$1.109\substack{+0.063\\-0.28}$	$1.053\substack{+0.079\\-0.21}$	$0.891\substack{+0.025\\-0.062}$	$0.893\substack{+0.021\\-0.065}$	$1.49\substack{+0.24 \\ -0.29}$	1.49 ± 0.26	$1.23\substack{+0.11\\-0.22}$	$1.15\substack{+0.10 \\ -0.14}$
$r_s \ [{ m Mpc}]$	c]	147.08 ± 0.30	147.12 ± 0.27	147.03 ± 0.25	147.05 ± 0.25	147.32 ± 0.27	147.35 ± 0.29	$147.31\substack{+0.25\\-0.29}$	$147.32\substack{+0.26 \\ -0.29}$
$\ln B_{ij}$		0.85	-0.17	1.60	0.60	-9.22	-11.68	-14.04	-15.21

A comparison between the **3D BAO data**, model dependent and obtained assuming ACDM, and the **2D BAO measurements**, less model dependent, shows completely different results for the IDE model. There is a strong evidence for the coupling at more than 99% CL, ⁷⁵ solving at the same time the H0 tension with SH0ES.

Bernui, Di Valentino, Giarè, Kumar, and Nunes, Phys. Rev. D 107 (2023) 10, 103531



Table II. Constraints at 68% CL on the parameters of the $\Lambda {\rm CDM}$ model.

Parameter	CMB	CMB+BAO-3D	CMB+BAO-2D (ON)	CMB+BAO-2D (M&M)
$10^2 imes \Omega_{ m b} h^2$	2.236 ± 0.015	2.245 ± 0.013	2.263 ± 0.014	2.246 ± 0.014
$\Omega_{ m c}h^2$	0.1202 ± 0.0014	0.11911 ± 0.00096	0.1165 ± 0.0011	0.11877 ± 0.00097
H_0	67.32 ± 0.62	67.84 ± 0.43	69.01 ± 0.51	67.96 ± 0.44
$ au_{ m reio}$	0.0536 ± 0.0081	0.0590 ± 0.0070	0.0606 ± 0.0081	0.0567 ± 0.0080
$\log(10^{10}A_{ m s})$	3.043 ± 0.016	3.053 ± 0.015	3.049 ± 0.017	3.047 ± 0.016
$n_{ m s}$	0.9646 ± 0.0045	0.9677 ± 0.0037	0.9742 ± 0.0038	0.9688 ± 0.0037

A comparison between the 3D BAO data and the 2D BAO measurements Menote & Marra arXiv:2112.10000, from the same BOSS DR12 and eBOSS DR16, gives exactly the same results for the ACDM scenario.

Table I. Constraints at 68% (95%) CL on the parameters of the IDE model.

Parameter	CMB	CMB+BAO-3D	CMB+BAO-2D (ON)	CMB+BAO-2D (M&M)
$10^2 imes \Omega_{ m b} h^2$	2.239 ± 0.015	2.236 ± 0.013	2.248 ± 0.014	2.237 ± 0.014
$\Omega_{ m c} h^2$	$0.067^{+0.042}_{-0.031} (< 0.115)$	$0.101\substack{+0.016\\ 0.012}$	$0.022\substack{+0.014\\-0.019}$	$0.089^{+0.019}_{-0.016}$
H_0	71.6 ± 2.1	$68.92\substack{+0.96\\-1.2}$	$75.2^{+1.1}_{-0.96}$	69.9 ± 1.1
$ au_{ m reio}$	0.0534 ± 0.0079	0.0544 ± 0.0079	0.0556 ± 0.0082	0.0537 ± 0.0078
$\log(10^{10}A_{ m s})$	3.042 ± 0.016	3.045 ± 0.016	3.044 ± 0.017	3.044 ± 0.016
$n_{ m s}$	0.9655 ± 0.0045	0.9650 ± 0.0037	0.9695 ± 0.0040	0.9657 ± 0.0039
ξ	$-0.40^{+0.23}_{-0.20} (> -0.775)$	> -0.207(> -0.389)	$-0.683\substack{+0.088\\-0.11}$	$-0.26^{+0.18}_{-0.12} (> -0.505)$

A comparison between the 3D BAO data and the 2D BAO measurements Menote & Marra arXiv:2112.10000, from the same BOSS DR12 and eBOSS DR16, gives different H0 values for the IDE scenario.

Parameter	Planck	ACT	ACT+WMAP	$\mathbf{ACT} + \mathbf{Planck}$
$\Omega_{ m b}h^2$	0.02237 ± 0.00015	0.02153 ± 0.00032	0.02238 ± 0.00020	0.02238 ± 0.00013
$\Omega_{ m c} h^2$	$0.067^{+0.042}_{-0.031} (< 0.115)$	< 0.0754 (< 0.111)	$0.070^{+0.046}_{-0.021} (< 0.117)$	$0.067^{+0.042}_{-0.030} (< 0.115)$
H_0	71.6 ± 2.1	$72.6\substack{+3.4 \\ -2.6}$	$71.3\substack{+2.6 \\ -3.2}$	$71.4^{+2.5}_{-2.8}$
$ au_{ m reio}$	0.0534 ± 0.0079	0.063 ± 0.015	0.061 ± 0.014	0.0533 ± 0.0073
$\log(10^{10}A_{ m s})$	3.042 ± 0.016	3.046 ± 0.030	3.064 ± 0.028	3.047 ± 0.014
$n_{ m s}$	0.9655 ± 0.0045	1.010 ± 0.016	$0.9741\substack{+0.0066\\-0.0064}$	0.9699 ± 0.0038
ξ	$-0.40\substack{+0.23\-0.20}$	$-0.46\substack{+0.20\\-0.28}$	$-0.38\substack{+0.35\\-0.14}$	$-0.40\substack{+0.27\\-0.23}$
$\ln \mathcal{B}_{ij}$	-0.17	-0.07	0.06	-0.25

Zhai, Giarè, van de Bruck, Di Valentino, et al, JCAP 07 (2023) 032

Let's now consider different combinations of CMB datasets.



Parameter	Planck	Planck ACT		$\mathbf{ACT} + \mathbf{Planck}$	
$\Omega_{ m b}h^2$	0.02237 ± 0.00015	0.02153 ± 0.00032	0.02238 ± 0.00020	0.02238 ± 0.00013	
$\Omega_{ m c} h^2$	$0.067^{+0.042}_{-0.031} (< 0.115)$	< 0.0754 (< 0.111)	$0.070^{+0.046}_{-0.021} (< 0.117)$	$0.067^{+0.042}_{-0.030} (< 0.115)$	
H_0	71.6 ± 2.1	$72.6\substack{+3.4 \\ -2.6}$	$71.3\substack{+2.6 \\ -3.2}$	$71.4^{+2.5}_{-2.8}$	
$ au_{ m reio}$	0.0534 ± 0.0079	0.063 ± 0.015	0.061 ± 0.014	0.0533 ± 0.0073	
$\log(10^{10}A_{ m s})$	3.042 ± 0.016	3.046 ± 0.030	3.064 ± 0.028	3.047 ± 0.014	
$n_{ m s}$	0.9655 ± 0.0045	1.010 ± 0.016	$0.9741\substack{+0.0066\\-0.0064}$	0.9699 ± 0.0038	
ξ	$-0.40\substack{+0.23\-0.20}$	$-0.46\substack{+0.20\\-0.28}$	$-0.38\substack{+0.35\\-0.14}$	$-0.40\substack{+0.27\\-0.23}$	
$\ln {\cal B}_{ij}$	-0.17	-0.07	0.06	-0.25	
	· · · · · · · · · · · · · · · · · · ·	· · · · · · · · · · · · · · · · · · ·			

Zhai, Giarè, van de Bruck, Di Valentino, et al, JCAP 07 (2023) 032

If we consider different combinations of CMB datasets, they provide similar results, favoring IDE with a 95% CL significance in the majority of the cases. Remarkably, such a preference remains consistent when cross-checked through independent probes, while always yielding a value of the expansion rate H0 consistent⁸¹ with the local distance ladder measurements.



Zhai, Giarè, van de Bruck, Di Valentino, et al, JCAP 07 (2023) 032

It is easy to observe that the preference for $\xi < 0$ is primarily driven by the high multipole ACT CMB data that have a reduced amplitude. These data are also responsible for the improvement of the fit in the context of IDE models compared to the minimal Λ CDM, indicating that it is a genuine effect rather than one caused by parameter degeneracies.

Let's see another example at late time...

Omnipotent DE

Density	EoS	Scaling in z	Scaling in a	Naming
	w>-1	$d\rho / dz > 0$	$\mathrm{d} ho /\mathrm{d}a < 0$	p-quintessence
ho > 0	w = -1	$\mathrm{d} ho/\mathrm{d}z=0$	$\mathrm{d} ho /\mathrm{d}a=0$	$\operatorname{positive-CC}$
	w < -1	$\mathrm{d} ho /\mathrm{d}z < 0$	$\mathrm{d} ho/\mathrm{d}a>0$	p-phantom
	w>-1	$d\rho/dz < 0$	$\mathrm{d} ho /\mathrm{d}a > 0$	n-quintessence
$\rho < 0$	w = -1	$\mathrm{d} ho /\mathrm{d}z = 0$	$\mathrm{d} ho /\mathrm{d}a=0$	negative-CC
	w < -1	$d\rho / dz > 0$	$\mathrm{d} ho/\mathrm{d}a < 0$	n-phantom

Adil, Akarsu, Di Valentino et al., arXiv:2306.08046

We named "Omnipotent DE" a class of phenomenologically DE models that are capable of incorporating all six combinations of negative and positive DE density (ρ_{DE} <0 and ρ_{DE} >0) with different equation of states w_{DE} <-1, w_{DE} =-1, and w_{DE} > -1 into a single expansion scenario for at least one point in its parameter space.

Phantom Crossing

A particular Omnipotent DE model is the one that introduces a transition in the dark energy density ρ_{DE} assuming that there is an extrema at a scale factor a_m. If we take a Taylor series expansion of ρ_{DE} around a_m, we have:

$$\rho_{DE}(a) = \rho_0 + \rho_2 (a - a_m)^2 + \rho_3 (a - a_m)^3$$

= $\rho_0 [1 + \alpha (a - a_m)^2 + \beta (a - a_m)^3].$

So the expansion rate of the Universe will be:

$$H^{2}(a)/H_{0}^{2} = \Omega_{m0}a^{-3} + \Omega_{k0}a^{-2} + \Omega_{\gamma0}a^{-4} + \left(\frac{1 - \Omega_{m0} - \Omega_{k0} - \Omega_{\gamma0}}{1 + \alpha(1 - a_{m})^{2} + \beta(1 - a_{m})^{3}}\right) \\ \left[1 + \alpha(a - a_{m})^{2} + \beta(a - a_{m})^{3}\right],$$

And the dark energy equation of state:

$$w_{DE}(a) = -1 - \frac{a[2\alpha(a-a_m) + 3\beta(a-a_m)^2]}{3[1 + \alpha(a-a_m)^2 + \beta(a-a_m)^3]}.$$

If $a_m < 1$, this crossing happens before the present day.

Phantom Crossing

Parameters	CMB+lensing	CMB+R19	CMB+BAO	CMB+Pantheon	CMB+all
a_m	< 0.276	> 0.830	0.859 ± 0.064	$0.917\substack{+0.054\\-0.029}$	$0.851^{+0.048}_{-0.031}$
lpha	< 17.7	< 8.62	7.3 ± 3.9	< 5.10	< 3.32
eta	< 16.7	16.0 ± 7.5	16.1 ± 7.8	$10.6^{+4.4}_{-7.9}$	$7.7^{+2.2}_{-4.7}$
$\Omega_c h^2$	0.1194 ± 0.0014	0.1196 ± 0.0014	0.1201 ± 0.0013	0.1198 ± 0.0014	0.1198 ± 0.0011
$\Omega_b h^2$	0.02243 ± 0.00014	0.02243 ± 0.00016	0.02238 ± 0.00014	0.02240 ± 0.00015	0.02240 ± 0.00014
$100 \theta_{MC}$	1.04097 ± 0.00031	1.04096 ± 0.00032	1.04092 ± 0.00030	1.04095 ± 0.00032	1.04093 ± 0.00030
au	0.0521 ± 0.0076	0.0532 ± 0.0080	$0.0539\substack{+0.0070\\-0.0080}$	0.0529 ± 0.0076	0.0521 ± 0.0075
n_s	0.9667 ± 0.0042	0.9665 ± 0.0045	0.9652 ± 0.0043	0.9659 ± 0.0045	0.9655 ± 0.0038
$\ln(10^{10}A_s)$	3.038 ± 0.015	3.041 ± 0.016	3.044 ± 0.016	3.041 ± 0.016	3.039 ± 0.015
$H_0[\mathrm{km/s/Mpc}]$	> 92.8	74.2 ± 1.4	$71.0^{+2.9}_{-3.8}$	$71.7^{+2.2}_{-3.1}$	70.25 ± 0.78
σ_8	$1.012^{+0.051}_{-0.009}$	0.881 ± 0.018	$0.848^{+0.027}_{-0.034}$	$0.860^{+0.026}_{-0.033}$	0.838 ± 0.011
S_8	$0.752^{+0.009}_{-0.025}$	0.818 ± 0.016	0.826 ± 0.019	0.828 ± 0.016	0.823 ± 0.011

We find that the combination of all the observational data including Planck, in agreement one with each other for this model, is indeed consistent with a_m < 1at more than 2σ.

Moreover this model also helps to alleviate the H0 tension between low and high redshift observations below 2σ , even for the full datasets combination.

Phantom Crossing

	Planck	Planck	Planck	Planck+BAO
Parameters	+BAO	+PantheonPlus	+ Panthe on Plus & SH0 ES	+PantheonPlus&SH0ES
$\Omega_{ m c}h^2$	0.1196 ± 0.0011	0.1193 ± 0.0013	0.1176 ± 0.0011	0.1198 ± 0.0011
$10^2 \Omega_{ m b} h^2$	2.241 ± 0.014	2.243 ± 0.015	2.257 ± 0.014	2.243 ± 0.014
$100 heta_{MC}$	1.04191 ± 0.00029	1.04193 ± 0.00029	1.04207 ± 0.00029	1.04191 ± 0.00029
au	0.0517 ± 0.0076	0.0516 ± 0.0077	$0.0490\substack{+0.0081\\-0.0071}$	0.0502 ± 0.0076
n_s	0.9662 ± 0.0038	0.9669 ± 0.0043	0.9714 ± 0.0041	0.9659 ± 0.0038
$\ln(10^{10}A_s)$	3.038 ± 0.015	3.037 ± 0.015	$3.027^{+0.016}_{-0.014}$	3.035 ± 0.015
a_m	$0.841\substack{+0.053\\-0.062}$	$0.952\substack{+0.042\\-0.016}$	$0.957^{+0.016}_{-0.023}$	$0.922\substack{+0.041\\-0.035}$
lpha	$7.2^{+4.8}_{-3.9}$	$3.4^{+1.5}_{-2.5}$	$7.0^{+1.6}_{-2.0}$	< 2.77
eta	> 11.1	$7.3^{+3.0}_{-6.0}$	$16.5^{+3.4}_{-4.3}$	$6.5^{+1.9}_{-3.4}$
$H_0[{ m km/s/Mpc}]$	$72.1^{+3.2}_{-4.5}$	$68.9^{+1.4}_{-2.0}$	73.49 ± 0.98	$\overline{70.05\pm0.64}$

The same results are confirmed for the new updated datasets BAO DR16 and Pantheon+, redeeming the possibility of a late time solution, if the DE is not monotonic and can be negative.

Phantom Crossing



Di Valentino et al., Entropy 23 (2021) 4, 404

Planck 2018, Astron.Astrophys. 641 (2020) A6

The CMB+BAO combination it is in better agreement with the phantom crossing than with the ΛCDM model.

Phantom Crossing



Adil, Akarsu, Di Valentino et al., arXiv:2306.08046

And the same it is true for the most updated full dataset combination Planck+BAO+PantheonPlus&SH0ES.

Phantom Crossing

	Phant	Phantom-crossing Dark Energy (PDE)					
Parameter	Base + LSS	$\frac{\text{Base}}{\text{Base}} + \text{LSS} + \text{S}_8$	$\frac{\text{Base}}{\text{Base}} + \text{LSS} + \text{S}_8 + \text{H}_0$				
a_m	$0.774_{-0.020}^{+0.037}$	$0.774\substack{+0.038\\-0.020}$	$0.773\substack{+0.044\\-0.025}$				
α	$8.2^{+2.6}_{-3.7}$	$8.0^{+2.5}_{-3.6}$	$6.5^{+1.6}_{-2.8}$				
β	$14.2^{+6.7}_{-8.7}$	$14.1\substack{+6.8 \\ -8.4}$	$10.2\substack{+4.5 \\ -6.9}$				
$100\omega_b$	2.246 ± 0.019	2.245 ± 0.018	2.246 ± 0.018				
$10\omega_{cdm}$	1.165 ± 0.015	1.166 ± 0.011	1.163 ± 0.010				
H_0	$75.70^{+2.05}_{-2.32}$	$75.60^{+1.93}_{-2.12}$	73.92 ± 1.09				
au	0.057 ± 0.005	0.057 ± 0.005	0.057 ± 0.005				
$\ln(10^{10}A_s)$	3.038 ± 0.012	3.038 ± 0.011	3.037 ± 0.011				
n_s	0.974 ± 0.006	0.974 ± 0.005	0.975 ± 0.005				
$r_{ m drag}$	$147. \pm 0.38$	147.91 ± 0.31	147.97 ± 0.30				
Ω_m	0.244 ± 0.015	0.245 ± 0.013	0.255 ± 0.007				
σ_8	0.854 ± 0.022	0.855 ± 0.021	0.839 ± 0.013				
S_8	0.770 ± 0.017	0.771 ± 0.010	0.774 ± 0.010				

Chudaykin et al., arXiv:2203.03666 [astro-ph.CO]

Base (SPT3G + Planck TT + Planck lensing) + LSS data prefer a phantom crossing at more than 4σ, and without including any priors is in agreement with H0 and S8.

Let's see another example of negative DE...

Sign-switching cosmological constant

The Λ_sCDM model is inspired by the recent conjecture that the universe went through a spontaneous AdS-dS transition characterized by a sign - switching cosmological constant:

$$\Lambda \quad
ightarrow \quad \Lambda_{
m s} \equiv \Lambda_{
m s0} \, {
m sgn}[z_{\dagger} - z],$$

Akarsu, E Di Valentino et al., arXiv:2307.10899

Data set	Planck	Planck+BAOtr	Planck+BAOtr	Planck+BAOtr	Planck+BAOtr
			+PP	+PP&SH0ES	+PP&SH0ES+KiDS-1000
Model	$\Lambda_{ m s}{ m CDM}$	$\Lambda_{ m s}{ m CDM}$	$\Lambda_{ m s}{ m CDM}$	$\Lambda_{ m s}{ m CDM}$	$\Lambda_{ m s}{ m CDM}$
	ACDM	ACDM	ACDM	ACDM	ACDM
z _†	unconstrained	$1.70^{+0.09}_{-0.19}(1.65)$	$1.87^{+0.13}_{-0.21}(1.75)$	$1.70^{+0.10}_{-0.13}(1.67)$	$1.72^{+0.09}_{-0.12}(1.70)$
$M_B[{ m mag}]$			$-19.317^{+0.021}_{-0.025}(-19.311)$	$-19.290\pm0.017(-19.278)$	$-19.282 \pm 0.017(-19.280)$
			$-19.407 \pm 0.013 (-19.411)$	$-19.379 \pm 0.012 (-19.373)$	$-19.372\pm0.011(-19.369)$
$H_0[{ m km/s/Mpc}]$	$70.77^{+0.79}_{-2.70}(71.22)$	$73.30^{+1.20}_{-1.00}(73.59)$	$71.72^{+0.73}_{-0.92}(71.97)$	$72.82 \pm 0.65(73.20)$	$73.16 \pm 0.64(73.36)$
	$67.39 \pm 0.55 (67.28)$	$68.84 \pm 0.48 (68.61)$	$68.55 \pm 0.44 (68.54)$	$69.57 \pm 0.42 (69.73)$	$69.83 \pm 0.37 (69.96)$
$\Omega_{ m m}$	$0.2860^{+0.0230}_{-0.0099}(0.2796)$	$0.2643^{+0.0072}_{-0.0090}(0.2618)$	$0.2768^{+0.0072}_{-0.0063}(0.2759)$	$0.2683 \pm 0.0052 (0.2646)$	$0.2646 \pm 0.0052(0.2622)$
	$0.3151 \pm 0.0075 (0.3163)$	$0.2958 \pm 0.0061 (0.2984)$	$0.2995 \pm 0.0056 (0.2992)$	$0.2869 \pm 0.0051 (0.2849)$	$0.2837 \pm 0.0045 (0.2816)$
S_8	$0.801^{+0.026}_{-0.016}(0.791)$	$0.777 \pm 0.011 (0.772)$	$0.791 \pm 0.011 (0.794)$	$0.783 \pm 0.010 (0.777)$	$0.774 \pm 0.009(0.773)$
	$0.832 \pm 0.013 (0.835)$	$0.802 \pm 0.011 (0.804)$	$0.808 \pm 0.010 (0.804)$	$0.788 \pm 0.010 (0.784)$	$0.781 \pm 0.008 (0.782)$
$\chi^2_{ m min}$	2778.06	2793.38	4219.68	4097.32	4185.34
	2780.52	2820.30	4235.18	4138.26	4226.50
$\ln \mathcal{B}_{ij}$	-1.28	-12.65	-7.52	-19.47	-19.77

Sign-switching cosmological constant

We see that there is no H₀ tension in the present analyses of Λ_sCDM with all data combinations including the BAOtr data, and this is is very strongly favored over ΛCDM in terms of Bayesian evidence. The favoured transition redshift is zt~1.7.

Data set	Planck	Planck+BAOtr	Planck+BAOtr	Planck+BAOtr	Planck+BAOtr
			+PP	+PP&SH0ES	+PP&SH0ES+KiDS-1000
Model	$\Lambda_{ m s}{ m CDM}$	$\Lambda_{ m s}{ m CDM}$	$\Lambda_{ m s}{ m CDM}$	$\Lambda_{ m s}{ m CDM}$	Λ_{s} CDM
	ACDM	ACDM	ACDM	ACDM	ACDM
<i>z</i> †	unconstrained	$1.70^{+0.09}_{-0.19}(1.65)$	$1.87^{+0.13}_{-0.21}(1.75)$	$1.70^{+0.10}_{-0.13}(1.67)$	$1.72^{+0.09}_{-0.12}(1.70)$
$M_B[{ m mag}]$			$-19.317^{+0.021}_{-0.025}(-19.311)$	$-19.290\pm0.017(-19.278)$	$-19.282\pm0.017(-19.280)$
			$-19.407\pm0.013(-19.411)$	$-19.379 \pm 0.012 (-19.373)$	$-19.372 \pm 0.011 (-19.369)$
$H_0[{ m km/s/Mpc}]$	$70.77^{+0.79}_{-2.70}(71.22)$	$73.30^{+1.20}_{-1.00}(73.59)$	$71.72^{+0.73}_{-0.92}(71.97)$	$72.82 \pm 0.65 (73.20)$	$73.16 \pm 0.64 (73.36)$
	$67.39 \pm 0.55 (67.28)$	$68.84 \pm 0.48 (68.61)$	$68.55 \pm 0.44 (68.54)$	$69.57 \pm 0.42 (69.73)$	$69.83 \pm 0.37 (69.96)$
$\Omega_{ m m}$	$0.2860^{+0.0230}_{-0.0099}(0.2796)$	$0.2643^{+0.0072}_{-0.0090}(0.2618)$	$0.2768^{+0.0072}_{-0.0063}(0.2759)$	$0.2683 \pm 0.0052 (0.2646)$	$0.2646 \pm 0.0052(0.2622)$
	$0.3151 \pm 0.0075 (0.3163)$	$0.2958 \pm 0.0061 (0.2984)$	$0.2995 \pm 0.0056 (0.2992)$	$0.2869 \pm 0.0051 (0.2849)$	$0.2837 \pm 0.0045 (0.2816)$
S_8	$0.801^{+0.026}_{-0.016}(0.791)$	$0.777 \pm 0.011(0.772)$	$0.791 \pm 0.011 (0.794)$	$0.783 \pm 0.010 (0.777)$	$0.774 \pm 0.009(0.773)$
	$0.832 \pm 0.013 (0.835)$	$0.802 \pm 0.011 (0.804)$	$0.808 \pm 0.010 (0.804)$	$0.788 \pm 0.010 (0.784)$	$0.781 \pm 0.008 (0.782)$
$\chi^2_{ m min}$	2778.06	2793.38	4219.68	4097.32	4185.34
	2780.52	2820.30	4235.18	4138.26	4226.50
$\ln \mathcal{B}_{ij}$	-1.28	-12.65	-7.52	-19.47	-19.77

Akarsu, E Di Valentino et al., arXiv:2307.10899

Sign-switching cosmological constant

We see that there is no H₀ tension in the present analyses of Λ_sCDM with all data combinations including the BAOtr data. Also the S8 tension is completely solved.



FIG. 2. 2D contours at 68%, and 95% CL in the H_0 - Ω_m plane for the Λ_s CDM and Λ CDM models from the Planck and/or BAOtr data. It deserves mention that, in case of Λ_s CDM, the Planck and BAOtr contours intersect right on the vertical band of SH0ES measurement.



Akarsu, E Di Valentino et al., arXiv:2307.10899

...but the excess of lensing in Planck could explain S8...

A_L internal anomaly

CMB photons emitted at recombination are deflected by the gravitational lensing effect of massive cosmic structures. The lensing amplitude AL parameterizes the rescaling of the lensing potential $\phi(n)$, then the power spectrum of the lensing field:

 $C_{\ell}^{\phi\phi} \to A_{\rm L} C_{\ell}^{\phi\phi}$

The gravitational lensing deflects the photon path by a quantity defined by the gradient of the lensing potential $\phi(n)$, integrated along the line of sight *n*, remapping the temperature field.

A_L internal anomaly

Its effect on the power spectrum is the smoothing of the acoustic peaks, increasing AL.

Interesting consistency checks is if the amplitude of the smoothing effect in the CMB power spectra matches the theoretical expectation AL = 1 and whether the amplitude of the smoothing is consistent with that measured by the lensing reconstruction.

If AL =1 then the theory is correct, otherwise we have a new physics or systematics.



Calabrese et al., Phys. Rev. D, 77, 123531

A_L : a failed consistency check

The Planck lensing-reconstruction power spectrum is consistent with the amplitude expected for Λ CDM models that fit the CMB spectra, so the Planck lensing measurement is compatible with AL = 1.

However, the distributions of AL inferred from the CMB power spectra alone indicate a preference for AL > 1.

The joint combined likelihood shifts the value preferred by the TT data downwards towards AL = 1, but the error also shrinks, increasing the significance of AL > 1 to 2.8σ .

The preference for high AL is not just a volume effect in the full parameter space, with the best fit improved by $\Delta\chi^2 \sim 9$ when adding AL for TT+lowE and 10 for TTTEEE+lowE.



Planck 2018, Astron.Astrophys. 641 (2020) A6

 $A_{\rm L} = 1.243 \pm 0.096$ (68 %, *Planck* TT+lowE), $A_{\rm L} = 1.180 \pm 0.065$ (68 %, *Planck* TT,TE,EE+lowE),

A_L can explain internal tension



A_L can explain internal tension



Marginalized 68.3% confidence Λ CDM parameter constraints from fits to the I < 1000 and I ≥ 1000 Planck TT 2015 spectra. Tension at more than 2 σ level appears in $\Omega_c h^2$ and derived parameters, including H0, Ω m, and σ 8.

Addison et al., Astrophys.J. 818 (2016) no.2, 132

A_L can explain internal tension



Increasing AL smooths out the high order acoustic peaks, improving the agreement between the two multipole ranges.

Addison et al., Astrophys.J. 818 (2016) no.2, 132

A_L can explain the S8 tension



Di Valentino, Melchiorri and Silk, JCAP 2001 (2020) no.01, 013

A_L that is larger than the expected value at about 3 standard deviations even when combining the Planck data with BAO and supernovae type Ia external datasets.

Alternative CMB are not in significant tension



CMB: Planck CMB aniso. CMB: Planck CMB aniso. (+A_{lens} marg.) CMB: WMAP+ACT CMB aniso. CMBL: Planck CMB lensing + BAO CMBL: SPT CMB lensing + BAO **CMBL: ACT CMB lensing + BAO CMBL: ACT +Planck CMB lensing + BAO** WL: DES-Y3 galaxy lensing+clustering WL: KiDS-1000 galaxy lensing+clustering HSC-Y3 galaxy lensing (Fourier) + BAO HSC-Y3 galaxy lensing (Real) + BAO GC: eBOSS BAO+RSD CX: SPT/Planck CMB lensing x DES CX: Planck CMB lensing x unWISE

ACT collaboration, arXiv:2304.05203



SPT-3G collaboration, arXiv:2212.05642

DES Y3 + KiDS-1000



DES Y3 + KiDS-1000 collaborations, arXiv:2305.17173 [astro-ph.CO]

There is no more S8 tension, showing now an agreement at about 1.7σ between Planck assuming ΛCDM and this combined analysis.

 $S_8 = 0.790^{+0.018} - 0.014$

DES Y3 + KiDS-1000 collaborations, arXiv:2305.17173 [astro-ph.CO]

But... assuming General Relativity, is there a physical explanation for A_L ?

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Planck 2018 results. VI. Cosmological parameters

Planck Collaboration: N. Aghanim⁵⁴, Y. Akrami^{15,57,59}, M. Ashdown^{65,5}, J. Aumont⁹⁵, C. Baccigalupi⁷⁸, M. Ballardini^{21,41}, A. J. Banday^{95,8},
 R. B. Barreiro⁶¹, N. Bartolo^{29,62}, S. Basak⁸⁵, R. Battye⁶⁴, K. Benabed^{55,90}, J.-P. Bernard^{95,8}, M. Bersanelli^{32,45}, P. Bielewicz^{75,78}, J. J. Bock^{63,10},
 J. R. Bond⁷, J. Borrill^{12,93}, F. R. Bouchet^{55,90}, F. Bou¹

J.-F. Cardoso^{55,90}, J. Carron²³, A. Challinor^{58,65,11}, H. C. Chia F. Cuttaia41, P. de Bernardis31, G. de Zotti42, J. Delabroui A. Ducout⁶⁶, X. Dupac³⁵, S. Dusini⁶², G. Efstathiou⁶⁵ J. Fergusson¹¹, R. Fernandez-Cobos⁶¹, F. Finelli^{41,47}, F. For S. Galli55,90[†], K. Ganga², R. T. Génova-Santos^{60,16}, M. A. Gruppuso^{41,47}, J. E. Gudmundsson^{94,25}, J. Hamann⁸⁶, Z. Huang⁸³, A. H. Jaffe⁵³, W. C. Jones²⁵, A. Karakci⁵⁹, N. Krachmalnicoff78, M. Kunz14,54,3, H. Kurki-Suonio24,4 M. Le Jeune², P. Lemos^{58,65}, J. Lesgourgues⁵⁶, F. Levri M. López-Caniego35, P. M. Lubin28, Y.-Z. Ma77,80,74, J. A. Marcos-Caballero⁶¹, M. Maris⁴³, P. G. Martin⁷, M. Mar P. R. Meinhold²⁸, A. Melchiorri^{31,50}, A. Mennella^{32,45} D. Molinari^{30,41,48}, L. Montier^{95,8}, G. Morgante⁴¹, A. Mo B. Partridge³⁹, G. Patanchon², H. V. Peiris²², F. Perrott J. P. Rachen¹⁸, M. Reinecke⁷², M. Remazeilles⁶⁴, A. B. Ruiz-Granados^{60,16}, L. Salvati⁵⁴, M. Sandri⁴¹, M. Savelair R. Sunyaev^{72,91}, A.-S. Suur-Uski^{24,40}, J. A. Tauber³⁶, D. Valenziano41, J. Valiviita24,40, B. Van Tent69, L. Vibert54. S. D. M. W (Affiliation

We present cosmological parameter results from the final isotropies, combining information from the temperature an improved measurements of large-scale polarization allow th cant gains in the precision of other correlated parameters. In many parameters, with residual modelling uncertainties estim spatially-flat 6-parameter ACDM cosmology having a power from polarization, temperature, and lensing, separately and i baryon density $\Omega_b h^2 = 0.0224 \pm 0.0001$, scalar spectral inde 68 % confidence regions on measured parameters and 95 % $100\theta_* = 1.0411 \pm 0.0003$. These results are only weakly dependent in many commonly considered extensions. Assuming the ba Hubble constant $H_0 = (67.4 \pm 0.5) \text{ km s}^{-1} \text{Mpc}^{-1}$; matter dens We find no compelling evidence for extensions to the base-A considering single-parameter extensions) we constrain the effective the Standard Model prediction $N_{\text{eff}} = 3.046$, and find that the to prefer higher lensing amplitudes than predicted in base AC from the ACDM model; however, this is not supported by BAO data. The joint constraint with BAO measurements on s with Type Ia supernovae (SNe), the dark-energy equation of constant. We find no evidence for deviations from a purely Keck Array data, we place a limit on the tensor-to-scalar ra deuterium abundances for the base-ACDM cosmology are in agreement with BAO, SNe, and some galaxy lensing observ including galaxy clustering (which prefers lower fluctuation measurements of the Hubble constant (which prefer a high favoured by the Planck data.

 $\Omega_K = -0.044^{+0.018}_{-0.015}$ (68%, *Planck* TT, TE, EE+lowE), (46b)

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a detection of curvature at about 3.4σ

an apparent detection of curvature at well over 2σ . The 99% probability region for the TT,TE,EE+lowE result is $-0.095 < \Omega_K < -0.007$, with only about 1/10000 samples at $\Omega_K \ge 0$. This is not entirely a volume effect, since the best-fit χ^2 changes by $\Delta \chi^2_{\text{eff}} = -11$ compared to base Λ CDM when adding the one additional curvature parameter. The reasons for the pull towards

*Corresponding author: G. Efstathiou, gpe@ast.cam.ac.uk [†]Corresponding author: S. Galli, gallis@iap.fr [‡]Corresponding author: A. Lewis, antony@cosmologist.info

*Corresponding author: G. Efstathiou gpe@est can ad uk Vature of the Universe [†]Corresponding author: S. Galli, gall (ap). [‡]Corresponding author: A. Lewis, antony@cosmologist.info

2019

Sep

Curvature of the universe

Can Planck provide an unbiased and reliable estimate of the curvature of the Universe? This may not be the case since a "geometrical degeneracy" is present with Ωm . When precise CMB measurements at arc-minute angular scales are included, since gravitational lensing depends on the matter density, its detection breaks the geometrical degeneracy. The Planck experiment with its improved angular resolution offers the unique opportunity of a precise measurement of curvature from a single CMB experiment. We simulated Planck, finding that such experiment could constrain curvature with a 2% uncertainty, without any significant bias towards closed models.



Di Valentino, Melchiorri and Silk, Nature Astron. 4 (2019) 2, 196-203

Curvature of the universe

Planck favours a closed Universe $(\Omega k < 0)$ with 99.985% probability. A closed Universe with $\Omega K = -0.0438$ provides a better fit to PL18 with respect to a flat model.

This is not entirely a volume effect, since the best-fit $\Delta \chi^2$ changes by -11 compared to base ACDM when adding the one additional curvature parameter. The improvement is due also to the fact that closed models could also lead to a large-scale cut-off in the primordial density fluctuations in agreement with the observed low

CMB anisotropy quadrupole.



Di Valentino, Melchiorri and Silk, Nature Astron. 4 (2019) 2, 196-203
Low CMB anisotropy quadrupole



Planck 2018, Astron.Astrophys. 641 (2020) A6

A model with $\Omega \kappa < 0$ is slightly preferred with respect to a flat model with AL > 1, because closed models better fit not only the damping tail, but also the lowmultipole data, especially the quadrupole.

Astrophysics

[Submitted on 5 Mar 2003 (v1), last revised 30 Jul 2003 (this version, v2)]

Is the Low CMB Quadrupole a Signature of Spatial Curvature?

G. Efstathiou (University of Cambridge)

The temperature anisotropy power spectrum measured with the Wilkinson Microwave Anisotropy Probe (WMAP) at high multipoles is in spectacular agreement with an inflationary Lambda-dominated cold dark matter cosmology. However, the low order multipoles (especially the quadrupole) have lower amplitudes than expected from this cosmology, indicating a need for new physics. Here we speculate that the low quadrupole amplitude is associated with spatial curvature. We show that positively curved models are consistent with the WMAP data and that the quadrupole amplitude can be reproduced if the primordial spectrum

truncates on scales comparable to the curvature scale.

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A lower quadrupole than predicted by the ACDM was already present in WMAP, and a closed universe to explain this effect was already taken into account.



Closed models predict substantially higher lensing amplitudes than in Λ CDM, because the dark matter content can be greater, leading to a larger lensing signal. The reasons for the pull towards negative values of Ω_K are essentially the same as those that lead to the preference for AL > 1.



Di Valentino, Melchiorri and Silk, Nature Astron. 4 (2019) 2, 196-203

A closed universe (Friedmann 1922) can explain AL!



Di Valentino, Melchiorri and Silk, Nature Astron. 4 (2019) 2, 196-203

A degeneracy between curvature and the AL parameter is clearly present. A closed universe can provide a robust physical explanation to the enhancement of the lensing amplitude. In fact, the curvature of the Universe is not new physics beyond the standard model, but it is predicted by the General Relativity, and depends on the energy content of the Universe.

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The evolution over time of the geometry of the universe is described by Einstein's equations:

$$R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = \frac{8\pi G}{c^2}T_{\mu\nu} + \Lambda g_{\mu\nu}$$

Adopting a 4-dimensional coordinate system for the space-time and the Cosmological Principle, i.e. a universe homogeneous and isotropic at large scales, the resulting metric is the Friedmann-Lemaitre-Robertson-Walker (FLRW), that describes the distance between two events in space-time.

$$ds^{2} = c^{2}dt^{2} - a^{2}(t)\left[\frac{dr^{2}}{1 - kr^{2}} + r^{2}\left(d\theta^{2} + \sin^{2}\theta d\varphi^{2}\right)\right]$$

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The curvature parameter k can be positive, null or negative, depending on the value of the curvature of the universe: positive, flat or negative.

$$R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = \frac{8\pi G}{c^2}T_{\mu\nu} + \Lambda g_{\mu\nu}$$

Adopting a 4-dimensional coordinate system for the space-time and the Cosmological Principle, i.e. a universe homogeneous and isotropic at large scales, the resulting metric is the Friedmann-Lemaitre-Robertson-Walker (FLRW), that describes the distance between two events in space-time.

$$ds^{2} = c^{2}dt^{2} - a^{2}(t) \left[\frac{dr^{2}}{1 - kr^{2}} + r^{2} \left(d\theta^{2} + \sin^{2}\theta d\varphi^{2} \right) \right]$$

Combining together the FLRW metric and Einstein's equations we obtain the Friedmann equations that describe the expansion history of the universe:

1st
$$H^2 = \left(\frac{\dot{a}}{a}\right)^2 = \frac{8\pi G}{3}\rho - \frac{k}{a^2} + \frac{\Lambda}{3}$$

2nd $\frac{\ddot{a}}{a} = -\frac{4\pi G}{3}\left(\rho + 3P\right) + \frac{\Lambda}{3}$

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If we divide the 1st Friedmann equation, for the critical density (density of a flat universe), we obtain today:

$$\Omega = \sum_{i} \Omega_{i} = \Omega_{m} + \Omega_{\Lambda} + \Omega_{r} = 1 - \Omega_{k}$$

From this equation it is possible to estimate the curvature of the universe, independently measuring the various contributions to the total density parameter Ω .

Figure: http://w3.phys.nthu.edu.tw

$$\begin{cases} \Omega > 1 \quad \Omega_k < 0 \\ \Omega = 1 \quad \Omega_k = 0 \\ \Omega < 1 \quad \Omega_k > 0 \end{cases} \xrightarrow{k>0} k > 0 \xrightarrow{k>0} : closed Universe \\ k = 0 : flat Universe \\ k < 0 : open Universe \end{cases}$$

What about Planck+BAO?



Planck 2018, Astron.Astrophys. 641 (2020) A6

Adding BAO data, a joint constraint is very consistent with a flat universe.

 $\Omega_K = 0.0007 \pm 0.0019$ (68 %, TT, TE, EE+lowE +lensing+BAO).

Given the significant change in the conclusions from Planck alone, it is reasonable to investigate whether they are actually consistent. In fact, a basic assumption for combining complementary datasets is that these ones must be consistent, i.e. they must plausibly arise from the same cosmological model. 120

BAO tension



Planck 2018, Astron. Astrophys. 641 (2020) A6

This is a plot of the acoustic-scale distance ratio, DV(z)/rdrag, as a function of redshift, taken from several recent BAO surveys, and divided by the mean acoustic-scale ratio obtained by Planck adopting a model. rdrag is the comoving size of the sound horizon at the baryon drag epoch, and DV, the dilation scale, is a combination of the Hubble parameter H(z) and the comoving angular diameter distance DM(z). 121

In a ACDM model the BAO data agree really well with the Planck measurements...

BAO tension



Di Valentino, Melchiorri and Silk, Nature Astron. 4 (2019) 2, 196-203

... but when we let curvature to vary

122

there is a striking disagreement between Planck spectra and BAO measurements!

BAO tension



In agreement with Handley, Phys.Rev.D 103 (2021) 4, L041301

What about Planck+FS?



The strong disagreement between Planck and BAO it is evident in this triangular plot, as well as that with the full-shape (FS) galaxy power spectrum measurements from the BOSS DR12 CMASS sample, at an effective redshift $z_{eff} = 0.57$.

Vagnozzi, Di Valentino, et al., Phys.Dark Univ. 33 (2021) 100851

$\Lambda CDM + \Omega_k$: a 7 parameter standard model

As it has been convincingly pointed out in Anselmi et al., arXiv:2207.06547, in absence of any theoretical arguments, we cannot use observations that suggest small Ω_k to enforce $\Omega_k=0$. The common practice of assuming $\Omega_k=0$ places the onus on proponents of "curved Λ CDM" to provide sufficient evidence that $\Omega_k \neq 0$, and this is required as an additional parameter. Given the current tensions in cosmological parameters and CMB anomalies this choice is at least open to debate. So it would be preferable to have the standard cosmological phenomenological model with at least 7 parameters.

Curvature can explain internal tension



Di Valentino, Melchiorri and Silk, Nature Astron. 4 (2019) 2, 196-203

In a closed Universe with $\Omega K = -0.045$, the cosmological parameters derived in the two different multipole ranges are now fully compatible.

Curvature can't explain external tensions



Di Valentino, Melchiorri and Silk, Nature Astron. 4 (2019) 2, 196-203

Varying $\Omega \kappa$, both the well known tensions on H0 and S8 are exacerbated. In a $\Lambda CDM + \Omega K$ model, Planck gives H0 = 54.4^{+3.3}-4.0 km/s/Mpc at 68% cl., increasing the tension with SH0ES at 5.5 σ , and S8 in disagreement at about 3.8 σ with KiDS-450, and more than 3.5 σ with DES.

What about non-CMB data?



Di Valentino, Melchiorri and Silk, Nature Astron. 4 (2019) 2, 196-203

It is now interesting to address the compatibility of Planck with combined datasets, like BAO + type-la supernovae + big bang nucleosynthesis data. In principle, each dataset prefers a closed universe, but BAO+SN-Ia+BBN gives H0 = 79.6 \pm 6.8 km/s/Mpc at 68%cl, perfectly consistent with SH0ES, but at 3.4 σ tension with Planck.

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BAO+SNIa+BBN+R18 gives sik = -0.091 ± 0.037 at 68%cl.

EFTOFLSS to investigate FS data

	$\ln(10^{10}\mathrm{A}_s)$	h	$\Omega_{cdm}h^2$	Ω_m	Ω_k	n_s	$2 * \log(\mathcal{L})$
Flat, fixed n_s	$2.85^{+0.11}_{-0.12}$ (3.03) $0.667^{+0.011}_{-0.011}$ (0.672)	$0.114^{+0.005}_{-0.004}$ (0.115)	$0.307^{+0.010}_{-0.011}$ (0.304)	-	-	367.2
Curved, fixed n_s	$2.55^{+0.21}_{-0.22}$ (2.77) $0.686^{+0.015}_{-0.016}$ (0.665)	$0.115^{+0.004}_{-0.005}$ (0.111)	$0.291^{+0.014}_{-0.013}$ (0.302)	$-0.089^{+0.049}_{-0.046}$ -0.042)	-	366.3
Flat, varying n_s	$2.80^{+0.14}_{-0.13}$ (2.97)) $0.669^{+0.012}_{-0.011}$ (0.668)	$0.117^{+0.009}_{-0.008}$ (0.114)	$0.312^{+0.017}_{-0.014}$ (0.304)	-	$0.950^{+0.04}_{-0.051}$ (0.972)	367.1
Curved, varying n_s	$2.19^{+0.29}_{-0.28}$ (2.62) $0.707^{+0.021}_{-0.021}$ (0.686)	$0.127^{+0.011}_{-0.009}$ (0.116)	$0.300^{+0.016}_{-0.014}$ (0.295	$-0.152^{+0.059}_{-0.053}(-0.089)$	$0.878^{+0.053}_{-0.055}$ (0.932)	364.8

Glanville et al., arXiv:2205.05892

0,10

In this paper they use EFTofLSS to simultaneously constrain measurements from the 6dFGS, BOSS, and eBOSS catalogues, in order to remove some of the assumptions of flatness that enter into other large-scale structure analyses. Fitting the FS data with a BBN prior they measure a >2σ preference for a closed universe.

Beyond six parameters: extending $\Lambda CDM + \Omega k$



Semenaite et al., arXiv:2210.07304

A similar result has been obtained by analysing a wKCDM model, and the parameter $\omega_{K=}\Omega_{k}h^{2}$ that gives

 $\omega_{\rm K} = -0.0116^{+0.0029}_{-0.0036}$

i.e. a 4σ preference for a closed universe.

Evidence for a phantom closed Universe at more than 99% CL!!



It is interesting to note that if a closed universe increases the fine-tuning of the theory, the removal of a cosmological constant reduces it. It is, therefore, difficult to decide whether a phantom closed model is less or more theoretically convoluted than ACDM.

Di Valentino, Melchiorri and Silk, ApJ Letters, 908, L9 (2021), arXiv:2003.04935

So... is the Universe closed?

What about Planck PR4 (NPIPE) with Camspec?

$ar \times iv > astro-ph > arXiv:2205.10869$



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Astrophysics > Cosmology and Nongalactic Astrophysics

[Submitted on 22 May 2022 (v1), last revised 11 Nov 2022 (this version, v2)]

CMB power spectra and cosmological parameters from Planck PR4 with CamSpec

Erik Rosenberg, Steven Gratton, George Efstathiou

We present angular power spectra and cosmological parameter constraints derived from the Planck PR4 (NPIPE) maps of the Cosmic Microwave Background. NPIPE, released by the Planck Collaboration in 2020, is a new processing pipeline for producing calibrated frequency maps from Planck data. We have created new versions of the CamSpec likelihood using these maps and applied them to constrain LCDM and single-parameter extensions. We find excellent consistency between NPIPE and the Planck 2018 maps at the parameter level, showing that the Planck cosmology is robust to substantial changes in the mapmaking. The lower noise of NPIPE leads to ~10% tighter constraints, and we see both smaller error bars and a shift toward the LCDM values for beyond-LCDM parameters including Omega_K and A_Lens.



PR4_12.6	A_L	Ω_K	$N_{ m eff}$	m_{ν}
TTTEEE	1.095 ± 0.056	$-0.025^{+0.013}_{-0.010}$	3.00 ± 0.21	< 0.161
TT	1.198 ± 0.084	$-0.042^{+0.022}_{-0.016}$	$2.98^{+0.28}_{-0.35}$	< 0.278
TE	0.96 ± 0.15	$-0.010^{+0.035}_{-0.015}$	$3.11_{-0.42}^{+0.38}$	< 0.400
EE	0.995 ± 0.15	$-0.012\substack{+0.034\\-0.017}$	4.6 ± 1.3	< 2.37
PR3_12.6	A_L	Ω_K	$N_{ m eff}$	m_{ν}
TTTEEE	1.146 ± 0.061	$-0.035^{+0.016}_{-0.012}$	$2.94^{+0.20}_{-0.23}$	< 0.143
TT	1.215 ± 0.089	$-0.047\substack{+0.024\\-0.017}$	$2.89^{+0.28}_{-0.32}$	< 0.248
TE	0.96 ± 0.17	$-0.015^{+0.043}_{-0.015}$	$2.96^{+0.42}_{-0.49}$	< 0.504
EE	1.15 ± 0.20	$-0.053\substack{+0.063\\-0.029}$	$2.46^{+0.94}_{-1.7}$	-

...but this new likelihood is not really solving the problem of AL/ΩK, that is mainly coming from the TT power spectrum. And the constraints coming from TT are not changing in the 2 releases...



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PR3_12.6 TTTEEE TT	A_L 1.146 ± 0.061 1.215 ± 0.089	Ω_K $-0.035^{+0.016}_{-0.012}$ $-0.047^{+0.024}_{-0.017}$	$\frac{N_{\rm eff}}{2.94^{+0.20}_{-0.23}}$ $2.89^{+0.28}_{-0.32}$	<i>m</i> _ν < 0.143 < 0.248
PR3_12.6 TTTEEE TT TE	A_L 1.146 ± 0.061 1.215 ± 0.089 0.96 ± 0.17	Ω_K $-0.035^{+0.016}_{-0.012}$ $-0.047^{+0.024}_{-0.017}$ $-0.015^{+0.043}_{-0.015}$	$\frac{N_{\rm eff}}{2.94^{+0.20}_{-0.23}}$ $2.89^{+0.28}_{-0.32}$ $2.96^{+0.42}_{-0.49}$	m_{ν} < 0.143 < 0.248 < 0.504

...but this new likelihood is not really solving the problem of AL/ΩK, that is mainly coming from the TT power spectrum.

And the constraints coming from TT are not changing in the 2 releases...

The constraints derived from the EE power spectrum are instead those pulling all the parameters towards ACDM and thus alleviating the tensions.



However, this change in EE is producing a significant shift of the acoustic scale parameter θ , and an internal tension at 2.8 σ between TT and EE, that becomes more than 3.2-3.3 σ when AL/ Ω K vary.

	ℓ range	N_D	$\hat{\chi}^2$	$(\hat{\chi}^2-1)/\sqrt{2/N_D}$
TT 143x143	30 - 2000	1971	1.021	0.67
TT 143x217	500 - 2500	2001	0.985	-0.47
TT 217x217	500 - 2500	2001	1.002	0.05
TT All	30 - 2500	5973	1.074	4.07
TE	30 - 2000	1971	1.055	1.73
EE	30 - 2000	1971	1.026	0.82
TEEE	20 - 2000	3942	1.046	2.02
TTTEEE	30 - 2500	9915	1.063	4.46

Table 1. χ^2 of the different components of the PR4_12.6 likelihood with respect to the TTTEEE best-fit model. N_D is the size of the data vector. $\hat{\chi}^2 = \chi^2/N_D$ is the reduced χ^2 . The last column gives the number of standard deviations of $\hat{\chi}^2$ from unity.

..but more significantly, the reduced χ^2 values show a more than 4σ tension of the data with the best-fit obtained by TTTEEE assuming a Λ CDM model.



A_L for different data releases

Table 1. Posterior A_L Constraints from Analyses of Planck Temperature and Polarization Data since 2018 Release

Reference	Data Version	Likelihood	Data Combination	A_L	' $N\sigma$ ' Preference
					for $A_L > 1$
Planck Collaboration VI (2020)	PR3/2018	plik	TTTEEE+lowl/lowE	1.180 ± 0.065	2.8σ
	PR3/2018	plik	TT+lowl/lowE	1.243 ± 0.096	2.5σ
Rosenberg et al. (2022)	PR3/2018	CamSpec	TTTEEE+lowl/lowE	1.146 ± 0.061	2.4σ
	PR3/2018	CamSpec	$\mathrm{TT}+\texttt{lowl/lowE}$	1.215 ± 0.089	2.4σ
	PR4/NPIPE	CamSpec	TTTEEE+lowl/lowE	1.095 ± 0.056	1.7σ
	PR4/NPIPE	CamSpec	TT+lowl/lowE	1.198 ± 0.084	2.4σ
Tristram et al. (2023)	PR4/NPIPE	HiLLiPoP	$\mathrm{TTTEEE}+\texttt{lowl/LoLLiPoP}^{\mathrm{a}}$	1.036 ± 0.051	0.7σ
	PR4/NPIPE	HiLLiPoP	$\mathrm{TT}+\texttt{lowl/LoLLiPoP}$	1.068 ± 0.081	0.8σ

Addison et al, arXiv:2310.03127

 $S_8 = 0.834 \pm 0.016$ H0 = 67.36 ± 0.54 km/s/Mpc

Planck 2018, Aghanim et al., arXiv:1807.06209 [astro-ph.CO]

 $S_8 = 0.819 \pm 0.014$ H0 = 67.64 ± 0.52 km/s/Mpc

Tristram et al., arXiv:2309.10034 [astro-ph.CO]

What about the alternative CMB experiments?

CMB Polarization Measurements with SPTpol

Nicholas Harrington UC Berkeley

$$\Omega_K = 0.001^{+0.018}_{-0.019}$$



SPT-3G, arXiv:2103.13618 [astro-ph.CO]



ACT-DR4 + WMAP gives at 68% CL

 $\Omega_k = -0.001 \pm 0.012$



ACT-DR4 2020, Aiola et al., arXiv:2007.07288 [astro-ph.CO]

Harrison-Zel'dovich scale-invariant spectrum?

Dataset	Scalar Spectral Index (n_s) ΛCDM
ACT	1.009 ± 0.015
ACT+BAO (DR12)	1.006 ± 0.013
ACT+BAO (DR16)	1.006 ± 0.014
ACT+DESy1	1.007 ± 0.013
ACT+SPT+BAO (DR12)	0.996 ± 0.012
Planck	0.9649 ± 0.0044
Planck+BAO (DR12)	0.9668 ± 0.0038
Planck+BAO (DR16)	0.9677 ± 0.0037
Planck $(2 \le \ell \le 650)$	0.9655 ± 0.0043
Planck ($\ell > 650$)	0.9634 ± 0.0085

Giarè, Renzi, Mena, Di Valentino, and Melchiorri, MNRAS 521 (2023) 2, 2911 ACT shows a preference for a larger spectral index consistent with a Harrison-Zel'dovich scale-invariant spectrum ns=1 of primordial density perturbations introducing a tension with a significance of 2.7σ with the results from the Planck satellite.
Harrison-Zel'dovich scale-invariant spectrum?



Harrison-Zel'dovich scale-invariant spectrum?

Dataset	Scalar Spectral Index (n_s) ACDM
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Giarè, Renzi, Mena, Di Valentino, and Melchiorri, MNRAS 521 (2023) 2, 2911

n ACT-DR4 2020, arXiv:2007.07288 [astro-ph.CO] this discrepancy was interpreted as a consequence of the lack of information concerning the first acoustic peak of the temperature power spectrum. To verify this origin of the discrepancy in the CMB values of ns, we have performed two separate analyses of the Planck observations, splitting the likelihood into low 2 < I < 650 and high I > 650 multipoles. We find that the discrepancy still persists at the level of 3σ (2σ) for low (high) multiple temperature data. Planck data still prefer a value of the scalar spectral index smaller than unity at $\sim 4.3\sigma$ when the information about the first acoustic peak is removed.



Giarè, Pan, Di Valentino, Yang, de Haro, and Melchiorri, arXiv:2305.15378

We tested some models of inflation regarded as well - established benchmark scenarios and found out that they are ruled out by ACT at more than 3σ.

In the plot we show for example the 2D contours at 68%, 95%, and 99% CL and 1D posteriors in the (n_s, N_{efolds}) plane for the Starobinsky model.
The grey vertical band refers to the typical range of folds expansion N_{efolds} ∈ [50, N_{max}], expected in standard inflation.
The upper limit, N_{max} ≤ 73, is represented by the black dashed line.

Very similar results are obtained for all the other potentials, and in particular for ACT we find the following values for the number of e-folds at 68% (95%) CL:

- $\mathcal{N} > 138$ ($\mathcal{N} > 92.8$) for the Starobinsky model;
- $\mathcal{N} > 134 \ (\mathcal{N} > 88.6)$ for α -Attractor models;
- $\mathcal{N} > 257 \ (\mathcal{N} > 208)$ for Polynomial inflation;
- $\mathcal{N} > 177 \ (\mathcal{N} > 105)$ for the SUSY potential.



Giarè, Renzi, Mena, Di Valentino, and Melchiorri, MNRAS 521 (2023) 2, 2911

Such preference remains robust under the addition of large scale structure information, and in the two-dimensional plane it can be definitely noted that the direction of the Ω_bh^2 - ns degeneracy is opposite for ACT and Planck, and the disagreement here is significantly exceeding 3σ . This tension is partially driven by the ACT polarization data, as we can see replacing it with the SPT polarization measurements, but while the tension is relaxed in the plane Ω_bh^2 - ns, this combination is still preferring ns=1.

Quantifying global CMB tension

Planck

Handley and Lemos, arXiv:2007.08496 [astro-ph.CC

ACT

SPT

0.024 $\Omega_b h^2$ 0.022 Dataset combination tension \boldsymbol{p} ACT vs *Planck* 0.86% 2.63σ ACT vs SPT 1.8% 2.37σ 0.120 $\Omega_c h^2$ *Planck* vs SPT 16.8% 1.38σ ACT vs Planck+SPT 0.52% 2.79σ 0.1051.044 $100 \theta_{MC}$ 1.0401.0360.120.080.04 $\ln(10^{10}A_s)$ 3.13.02.91.02 n_s 0.96 $0.105 \ 0.120 \ 1.036 \ 1.040 \ 1.044 \ 0.04 \ 0.08 \ 0.122.9$ 3.03.10.960.0220.0241.02 $\Omega_b h^2$ $\Omega_c h^2$ $100\theta_{MC}$ $\ln(10^{10}A_{s})$ au n_{s}

Global tensions between CMB datasets.

For each pairing of datasets this is the tension probability p that such datasets would be this discordant by (Bayesian) chance, as well as a conversion into a Gaussianequivalent tension.

Between Planck and ACT there is a 2.6σ tension.

Assuming ACD/M

ACT-DR4 vs Planck: EDE



ACT collaboration, Hill et al. arXiv:2109.04451

Considering ACT only data or combined with Planck TT up to multipoles 650, there is an evidence for EDE > 3σ, solving completely the Hubble tension. The evidence for EDE > 3σ persists with the inclusion of Planck lensing + BAO data, but shifting H0 towards a lower value. Once the full Planck data are considered, the evidence for EDE disappears and H0 is again in tension with SH0ES.

The Planck damping tail is in disagreement with EDE different from zero.

ACT-DR4 vs Planck: α_s and β_s

Forconi, Giarè, Di Valentino and Melchiorri, *Phys.Rev.D* 104 (2021) 10, 103528



ACT-DR4 and SPT-3G are in agreement one with each other, but in disagreement with Planck, for the value of the

running of the scalar spectral index α_s and of the running of the running β_s . In particular ACT-DR4 + WMAP prefer both a non vanishing running α_s and running of the running β_s at the level of 2.9 σ and 2.7 σ , respectively.

Alternative CMB vs Planck: Σmv



Planck 2018 collaboration, arXiv:1807.06209 [astro-ph.CO]

While we have only an upper limit for Planck on the total neutrino mass, ACT-DR4, when combined with WMAP and lensing, prefers a neutrino mass different from zero at more than 95% CL.



Di Valentino and Melchiorri, 2022 ApJL 931 L18

Constraints at 68% CL					
Dataset	$\Sigma m_{\nu} [\mathrm{eV}]$				
ACT-DR4+WMAP+Lensing	0.60 ± 0.25				
${ m Planck+Lensing}\;(+A_{ m lens})$	$0.41\substack{+0.17 \\ -0.25}$				

Quantifying global CMB tension

Cosmological model	d	χ^2	p	$\log S$	Tension
ΛCDM	6	16.3	0.012	-5.17	2.51σ
ACDM + 4	7	18.5	0.00977	_5 77	258σ
$\Lambda { m CDM} + N_{ m eff}$	7	13	0.0719	-3	1.80σ
$\Lambda ext{CDM} + \Omega_k$	7	16.5	0.0209	-4.75	2.31σ
$w ext{CDM}$	7	16.8	0.0187	-4.9	2.35σ
$\Lambda ext{CDM} + \sum m_{ u}$	7	20.7	0.00421	-6.86	2.86σ
$\Lambda \text{CDM} + \alpha_s$	7	20.6	0.00448	-6.78	2.84σ
$w ext{CDM} + \Omega_k$	8	17.6	0.0249	-4.78	2.24σ
$\Lambda ext{CDM} + \Omega_k + \sum m_ u$	8	21.2	0.00651	-6.62	2.72σ
$w ext{CDM} + \Omega_k + \sum m_{ u}$	9	19.8	0.0195	-5.38	2.34σ
$w{ m CDM}+\Omega_k+\sum m_ u+N_{ m eff}$	10	18.8	0.0434	-4.38	2.02σ
$w \text{CDM} + \Omega_k + \sum m_\nu + \alpha_s$	10	22	0.015	-0.01	2.43σ
$w \text{CDM} + \Omega_k + N_{\text{eff}} + \alpha_s$	10	20.9	0.0218	-5.45	2.29σ
w CDM + $\sum m_{\nu} + N_{\text{eff}} + \alpha_s$	10	31.1	0.000575	-10.5	3.44σ
$w ext{CDM} + \Omega_k + \sum m_ u + N_{ ext{eff}} + lpha_s$	11	24.7	0.0102	-6.83	2.57σ

Di Valentino et al., MNRAS 520 (2023) 1, 210-215

$\Lambda { m CDM} + N_{ m eff}$	Planck	_	2.92 ± 0.19
	ACT-DR4	_	$2.35\substack{+0.40 \\ -0.47}$

If we now study the global agreement between Planck and ACT in various cosmological models that differ by the inclusion of different combinations of additional parameters, we can use the Suspiciousness statistic, to quantify their global "CMB tension".

We find that the 2.5 σ tension within the baseline Λ CDM is reduced at the level of 1.8 σ when Neff is significantly less than 3.044, while it ranges between 2.3 σ and 3.5 σ in all the other extended models.

Concluding

At this point, given the quality of all the analyses at play, probably these tensions are indicating a problem with the underlying cosmology and our understanding of the Universe, rather than the presence of systematic effects.

Many models have been proposed to solve the H0 tension. However, looking for a solution by changing the standard model of cosmology is challenging because of some additional complications:

- 1. The sound horizon problem
- 2. The S8 tension
- 3. The correlation between the parameters and possible fake detection
- 4. The hidden model dependence of some of the datasets (such as BAO)
- 5. The Planck AL problem
- 6. The inconsistency between the different CMB experiments

Therefore, this is presenting a serious limitation to the precision cosmology.

These cosmic discordances

call for new observations and stimulate the investigation of alternative theoretical models and solutions.



Thank you! e.divalentino@sheffield.ac.uk

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Addressing observational tensions in cosmology with systematics and fundamental physics

https://cosmoversetensions.eu/

WG1 – Observational Cosmology and systematics

Unveiling the nature of the existing cosmological tensions and other possible anomalies discovered in the future will require a multi-path approach involving a wide range of cosmological probes, various multiwavelength observations and diverse strategies for data analysis.

WG2 – Data Analysis in Cosmology

Presently, cosmological models are largely tested by using well-established methods, such as Bayesian approaches, that are usually combined with Monte Carlo Markov Chain (MCMC) methods as a standard tool to provide parameter constraints.

WG3 - Fundamental Physics

Given the observational tensions among different data sets, and the unknown quantities on which the model is based, alternative scenarios should be considered.



